

High brightness drive beam generation and kinetic plasma instabilities relevant to acceleration in dense plasmas

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- **Drivers needed for exciting wakes in dense plasmas (crystals & nano-structures)**

- **X-ray**

- optical laser (surface compression)
- e- beam (FEL)

- **10 keV?**

- **Electron beam**

- optical laser (LWFA)
- e- beam (PWFA)

- **10s MeV-GeV**

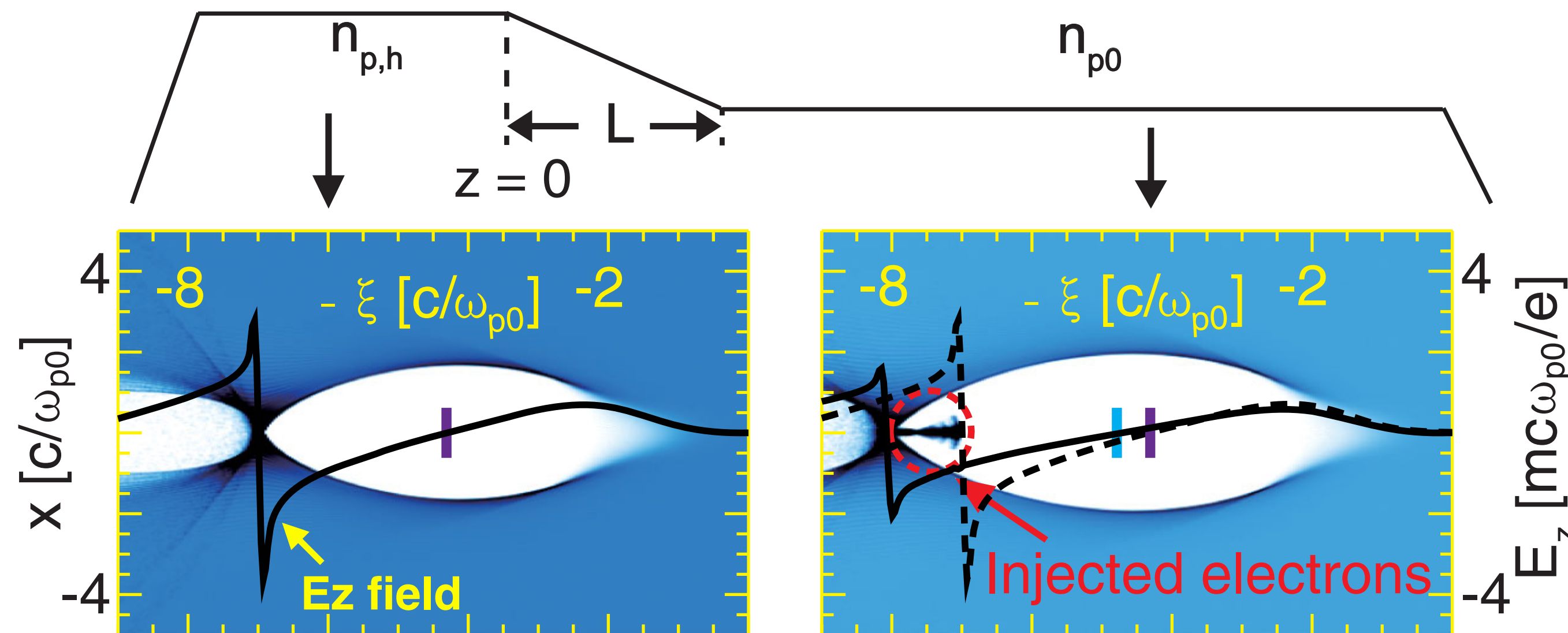
- **high brightness**

- **This talk will cover two topics:**

- **generation of high brightness e- beams in a PWFA using downramp trapping**
- **kinetic plasma instabilities relevant to beam-plasma interaction**

- **High brightness beam generation using downramp trapping**
 - **concept**
 - **simulation results**
- **Kinetic instabilities relevant to acceleration in dense plasmas**
 - **kinetic instabilities in beam-plasma system**
 - **experimental investigation using optical-field ionized plasmas**
 - **current filamentation experiment at FACET-II**

- Conceptual illustration of downramp trapping:



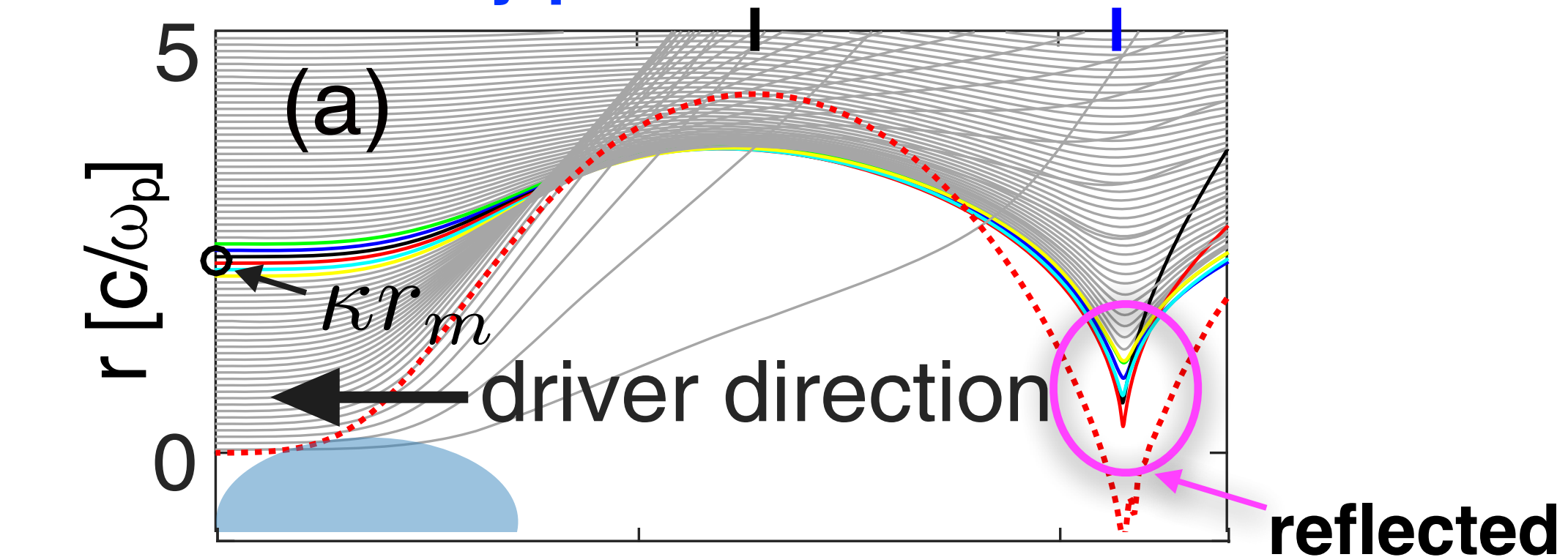
Xinlu Xu et al., Phys. Rev. Accel. Beams (2017)

The concept of downramp trapping was firstly proposed by S. Bulanov (1998) and H. Suk (2003) using 1D analysis.

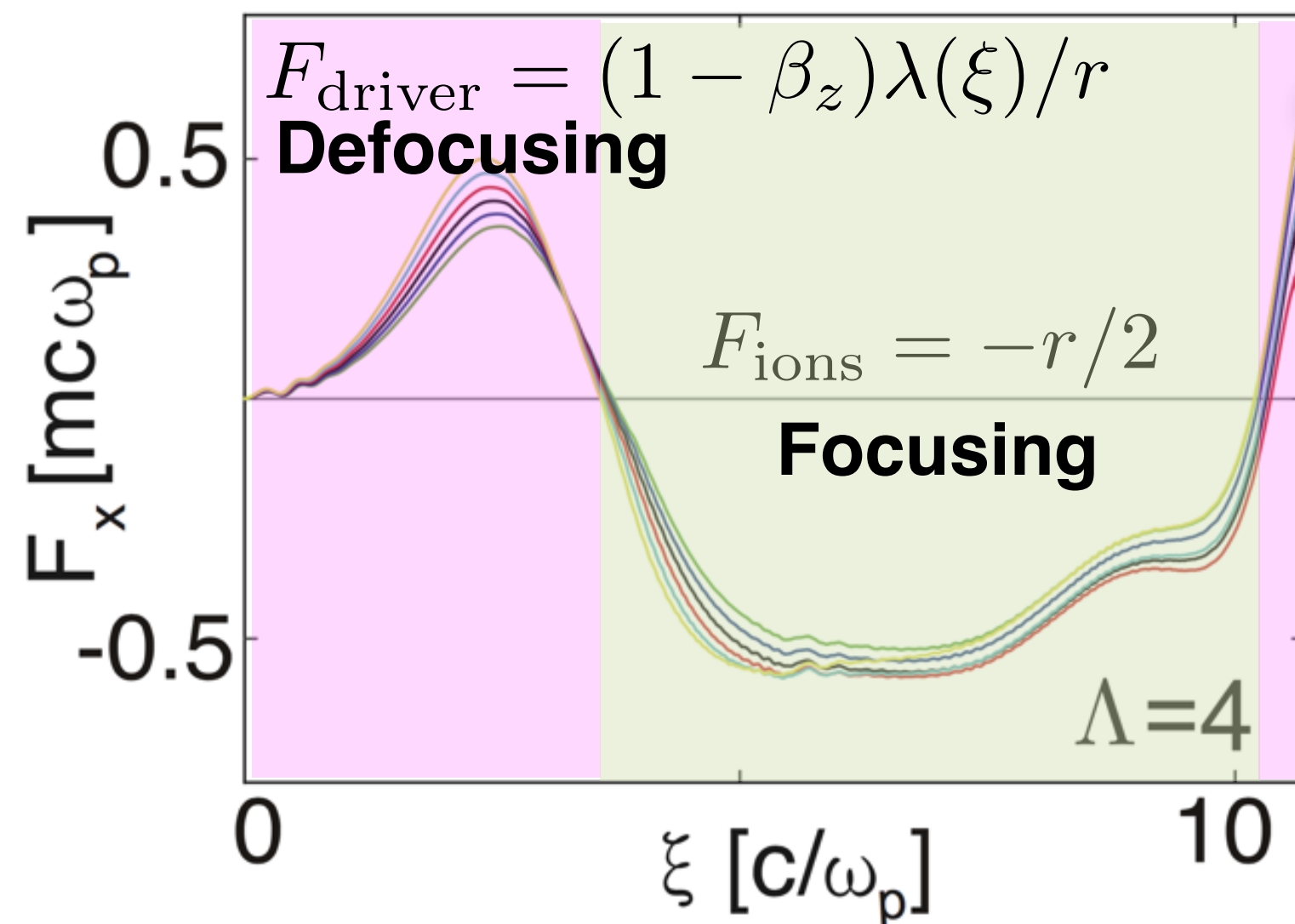
S. Bulanov, et al., Phys. Rev. E 58, R5257 (1998);

H. Suk, et al., Phys. Rev. Lett. 86, 1011 (2001)

Trajectories of electrons forming a wake in an uniform density plasma:



- as the electrons being pulled back to the axis, their radial velocity initially increases, but is reduced as these electrons approach the axis (actually most sheath electrons are reflected to form a second bubble)

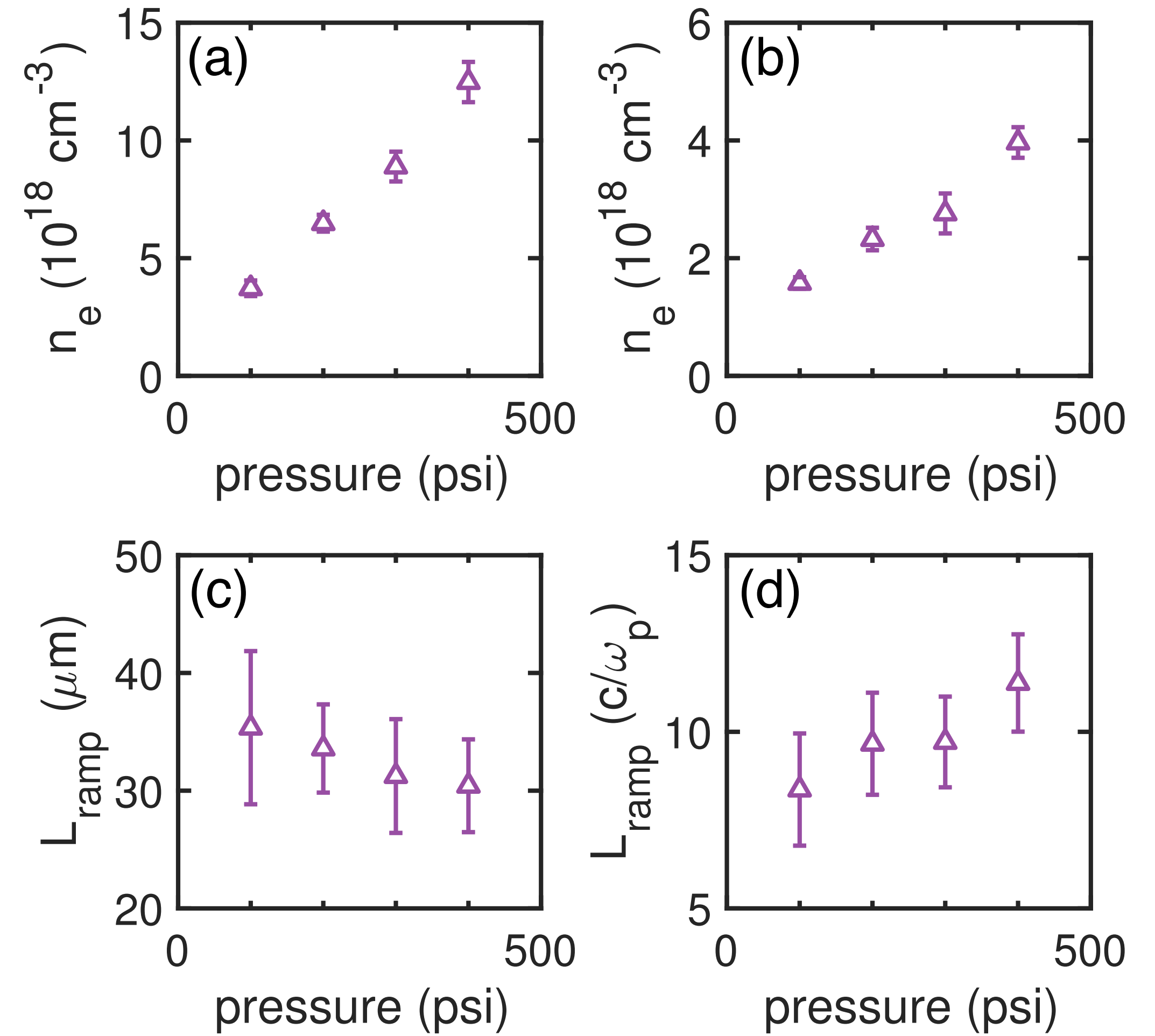
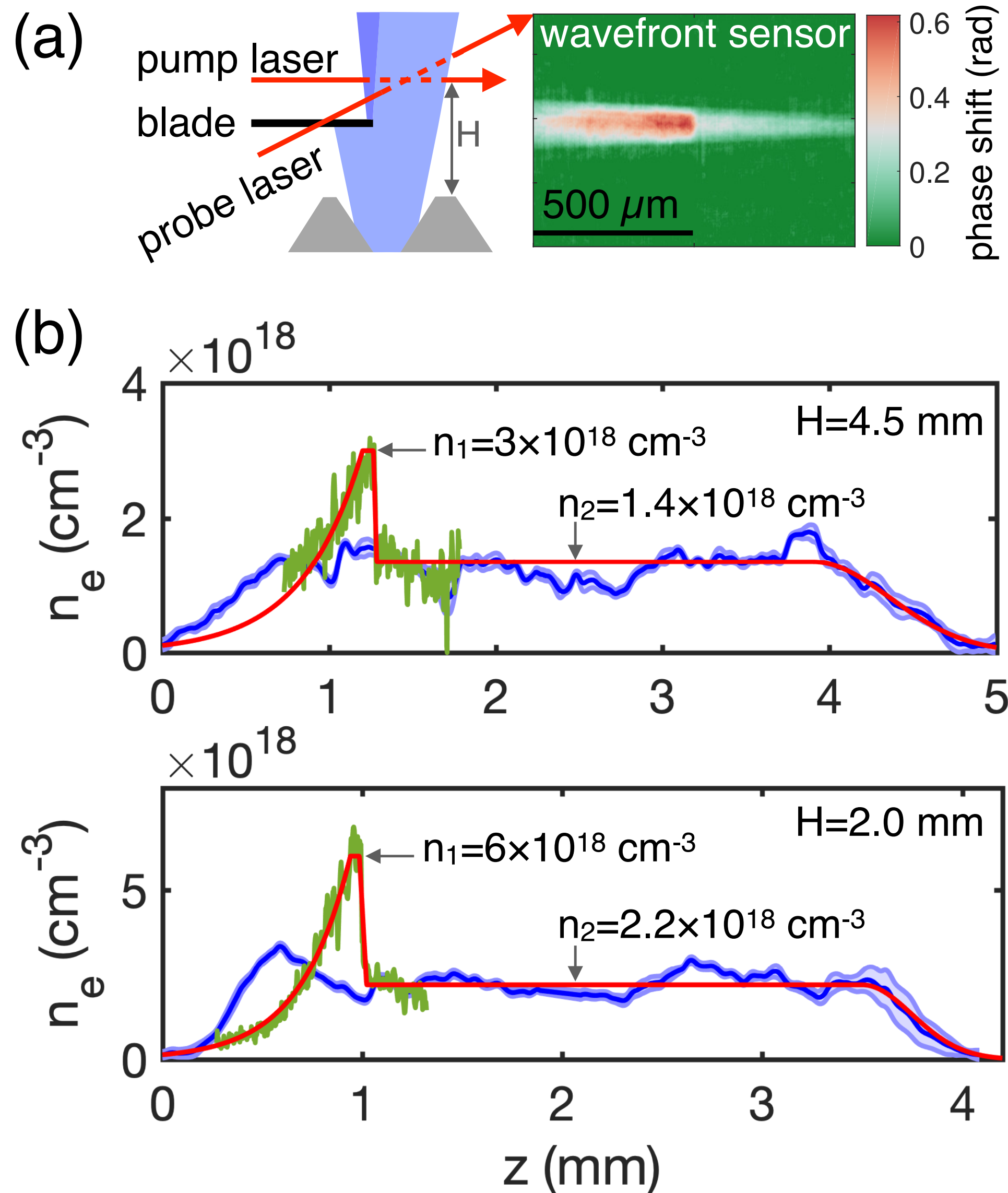


$$F_e = -(r/2)(1 - \beta_z)d^2\psi_0^2/d\xi^2$$

Defocusing

$$|\frac{d\psi_0^2}{d\xi^2}| \gg 1$$

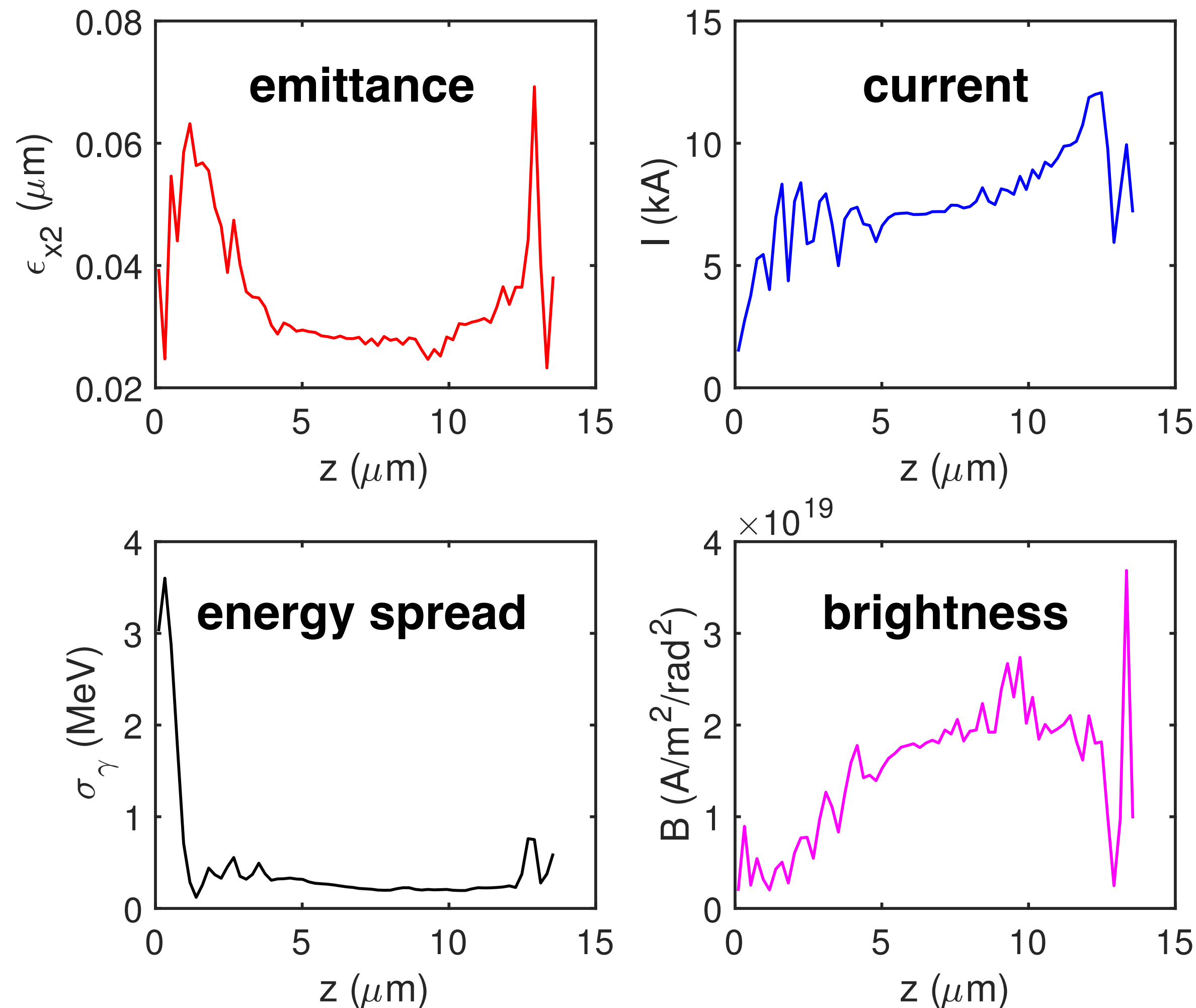
- beam current $\Lambda > 1$
- ramp length $\gg c/\omega_p$



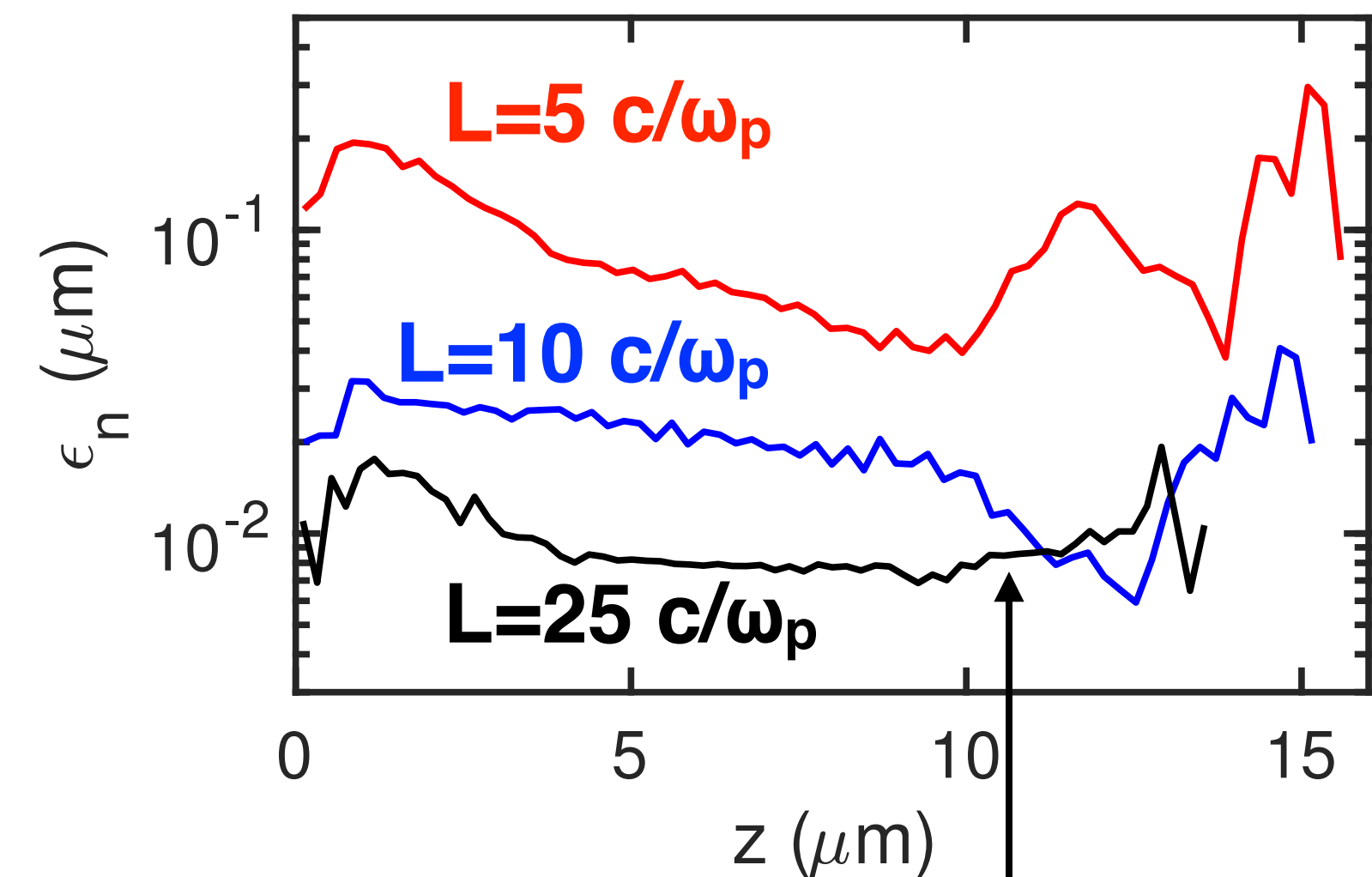
ramp length $\sim 10\ c/\omega_p$, adiabatic injection

A simulation example to show the parameters of the injected bunch:

downramp: $6e18 \text{ cm}^{-3}$ to $2.2e18 \text{ cm}^{-3}$ in $L=25 \text{ c}/\omega_p$



ϵ_n as a function of ramp length:



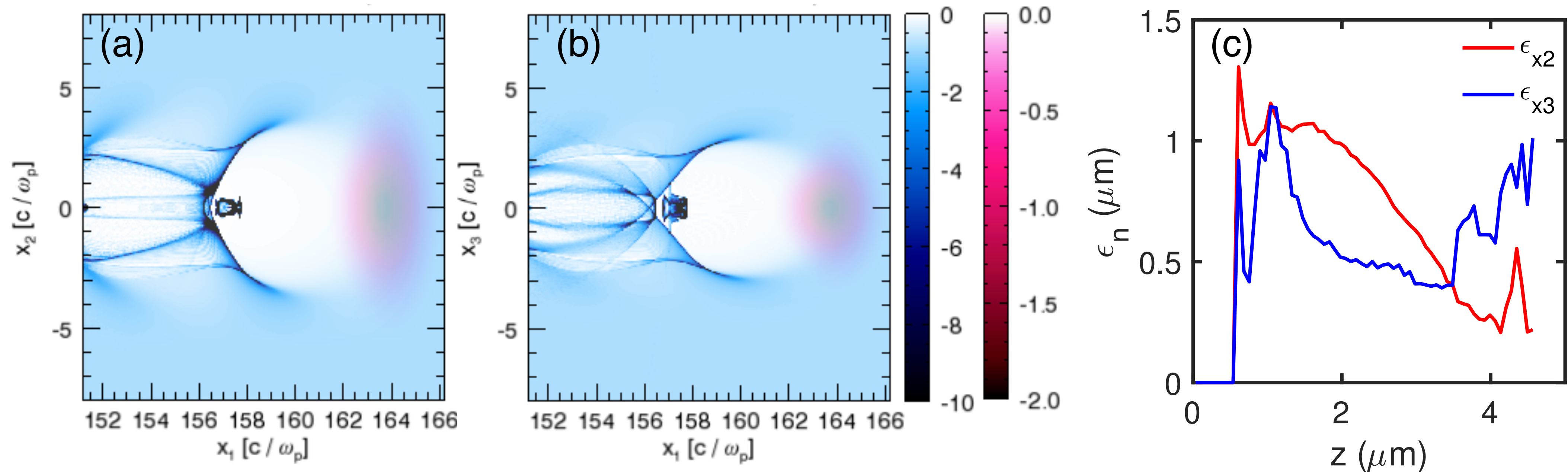
10 nm which is about 2 orders of magnitude lower than the current available values in plasma-based accelerators.

The emittance is sensitive to the symmetry of the wake (driver)

downramp: $3e18 \text{ cm}^{-3}$ to $2.2e18 \text{ cm}^{-3}$ in $L=10 \text{ c}/\omega_p$

$\sigma_{x2}=6.6 \text{ } \mu\text{m}$, $\sigma_{x3}=4.5 \text{ } \mu\text{m}$

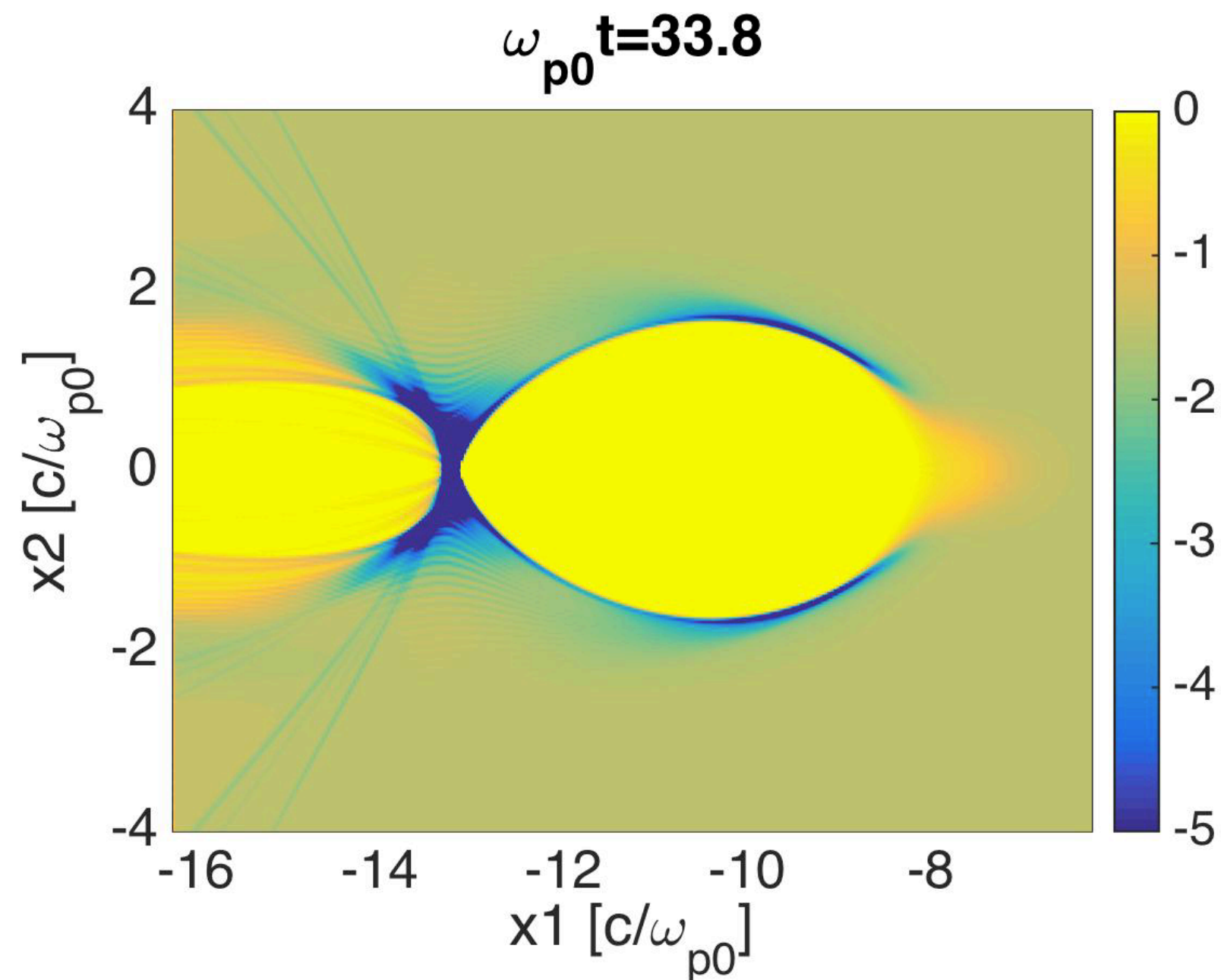
the slice emittance increases by a factor of ~ 10 compared to the symmetric driver case



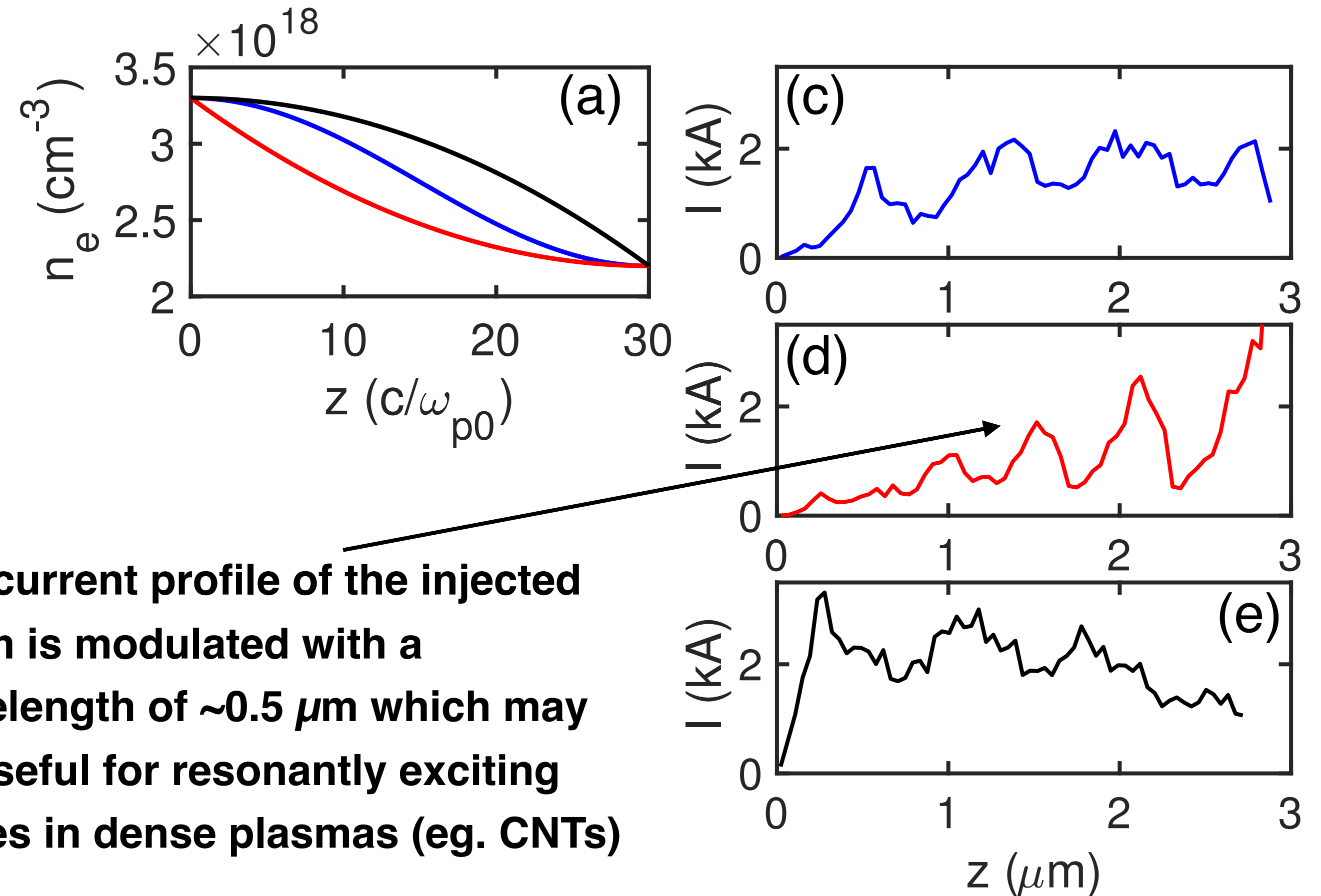
The current profile of the bunch depends on ramp shape

there is a one-to-one mapping between the initial (z_i) and the final (ξ_f) coordinates of the injected electrons, which will

- introduce an initial energy chirp
- affect the current profile



current profile depends on ramp shape



The current profile of the injected beam is modulated with a wavelength of $\sim 0.5 \mu\text{m}$ which may be useful for resonantly exciting wakes in dense plasmas (eg. CNTs)

The emittance and brightness of the injected beam scale with plasma density

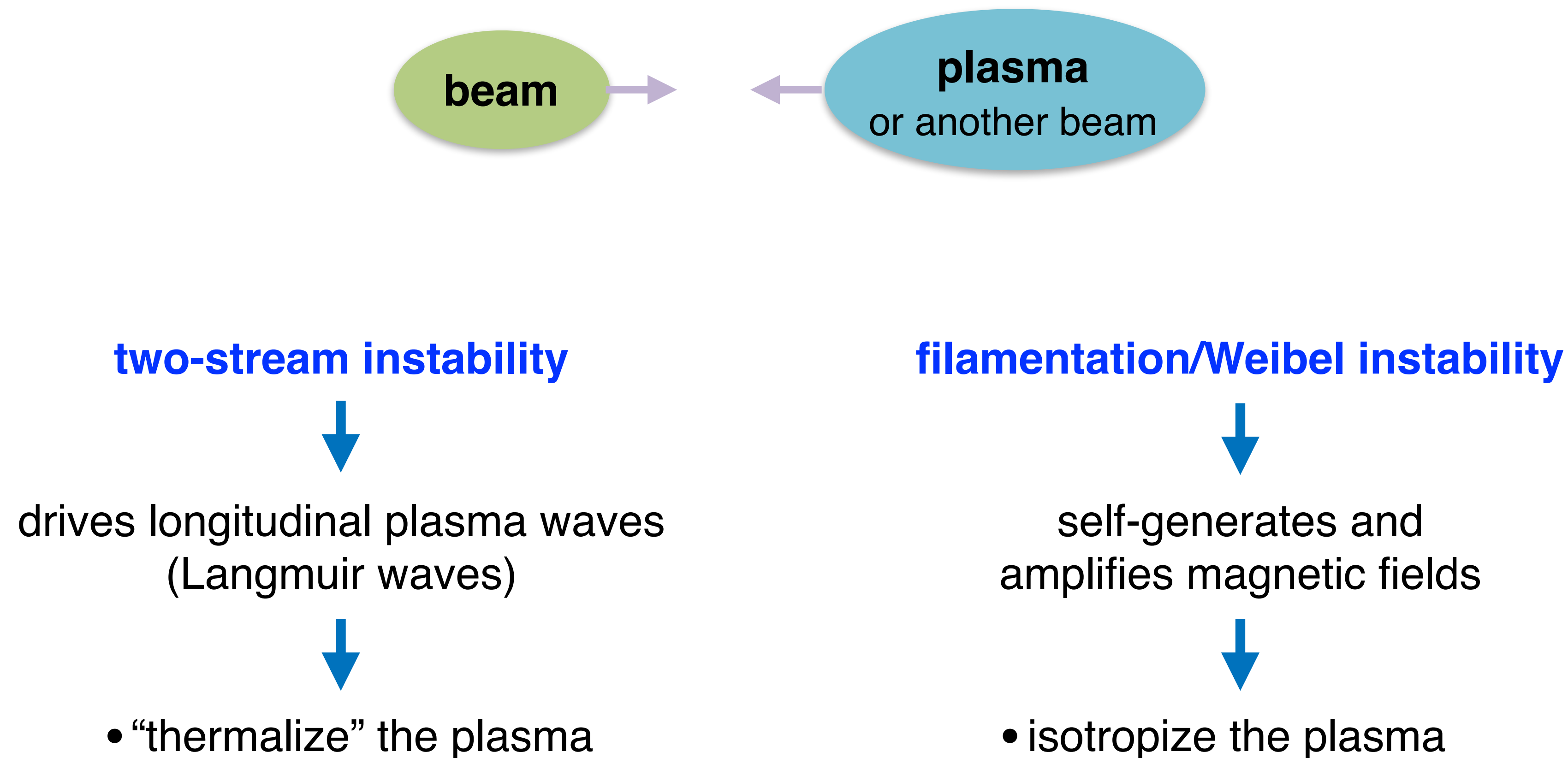
| | I | $\epsilon_n, \tau, \sigma_r$ | B | σ_E/E | n | Q |
|---------|-------------------------|------------------------------|------------|-----------------------------|------------|-----------------|
| E beams | n_{p0}^0 | $n_{p0}^{-0.5}$ | n_{p0}^1 | n_{p0}^0 | n_{p0}^1 | $n_{p0}^{-0.5}$ |
| | Ramp length | | | Optimal acceleration length | | |
| Plasma | $\propto n_{p0}^{-0.5}$ | | | $\propto n_{p0}^{-0.5}$ | | |

Case study: 1.5 $n_{p0} \rightarrow n_{p0}$, ramp length $L=250 \text{ c}/\omega_{p0}$:

| | n_{p0} [cm ⁻³] | I [kA] | ϵ_n [nm] | τ [fs] | σ_r [μm] | B [A/m ² /rad ²] | E [MeV] | σ_E/E | Q [pC] |
|---------------|---------------------------------|------------------------------------|----------------------|------------------------|--------------------|--|------------|----------------------|-----------|
| Injected beam | 10^{18} | 14 | 80 | 10 | 0.2 | 4×10^{18} | 620 | 1.5×10^{-3} | 140 |
| | 10^{20} | 14 | 8 | 1 | 0.02 | 4×10^{20} | 620 | 1.5×10^{-3} | 14 |
| | n_{p0} [cm ⁻³] | Density change [cm ⁻³] | | L_{ramp} [mm] | | L_{acc} [mm] | | | |
| Plasma | 10^{18} | 5×10^{17} | | 1.33 | | 3.3 | | | |
| | 10^{20} | 5×10^{19} | | 0.133 | | 0.33 | | | |

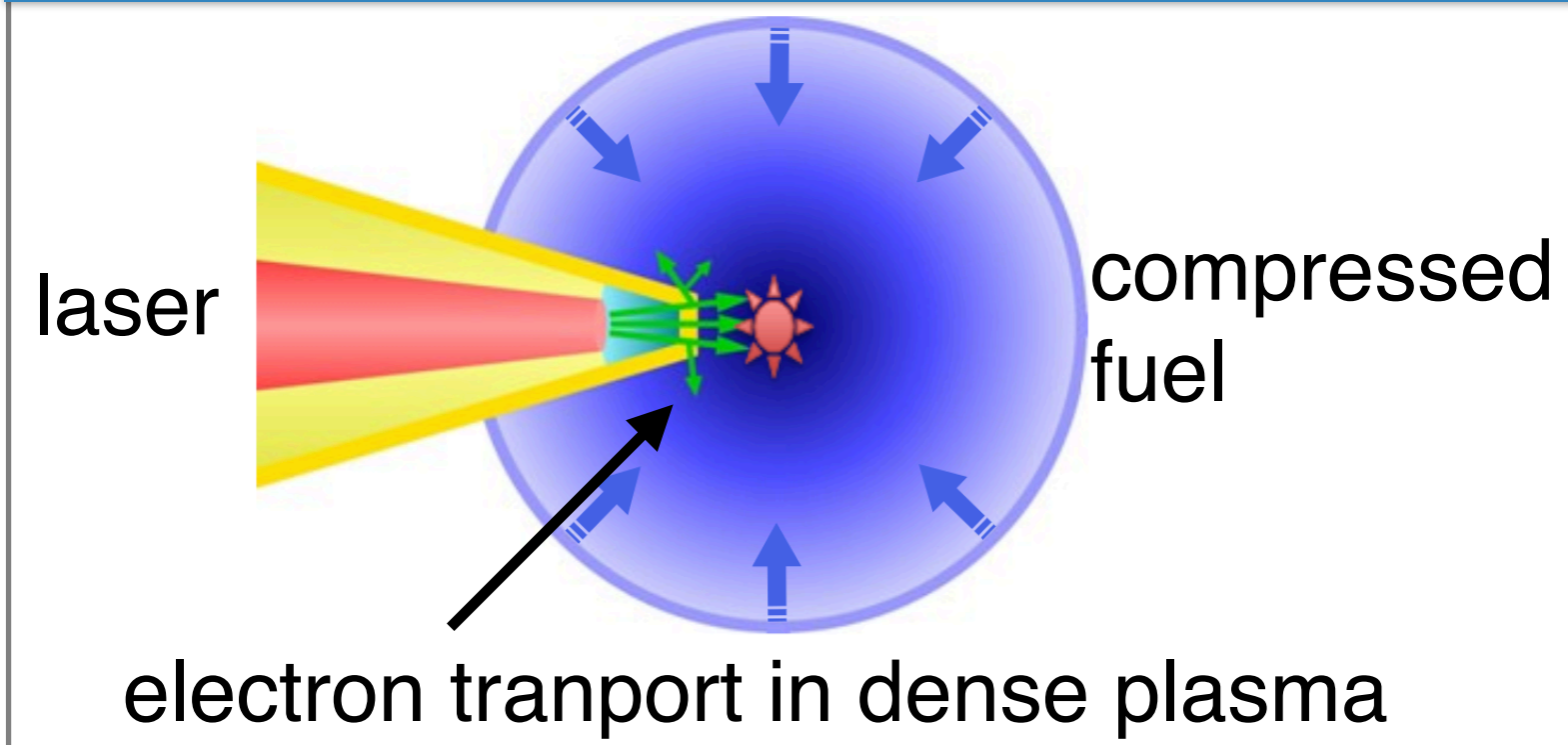
- High brightness beam generation using downramp trapping
 - concept
 - simulation results
 - downramp trapping experiment at FACET-II
- **Kinetic instabilities relevant to acceleration in dense plasmas**
 - **kinetic instabilities in beam-plasma system**
 - **experimental investigation using optical-field ionized plasmas**
 - **current filamentation experiment at FACET-II**

A beam-plasma system may be unstable to **two-stream** and current **filamentation** instability.

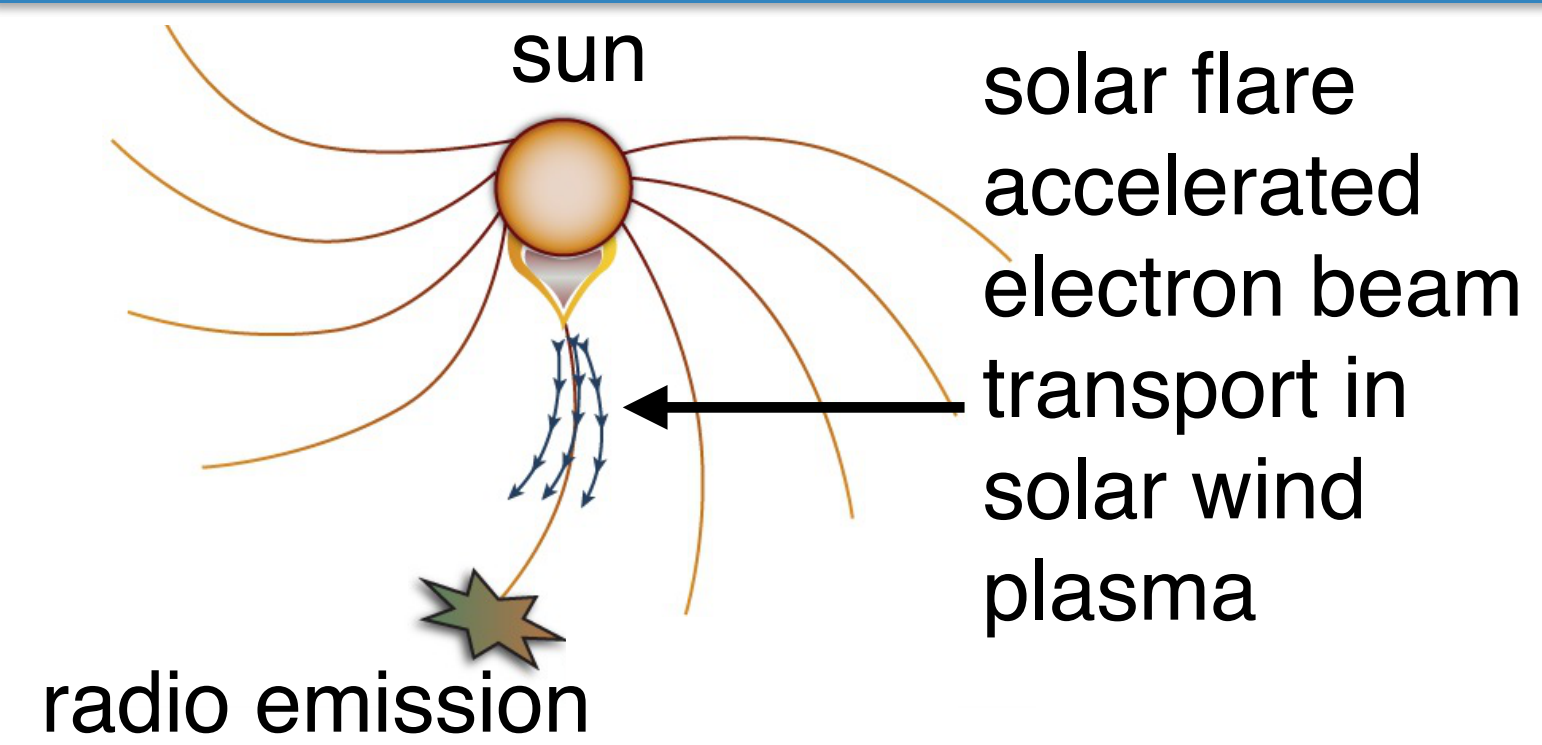


Kinetic plasma instabilities play important roles in both laboratory and astrophysical plasmas

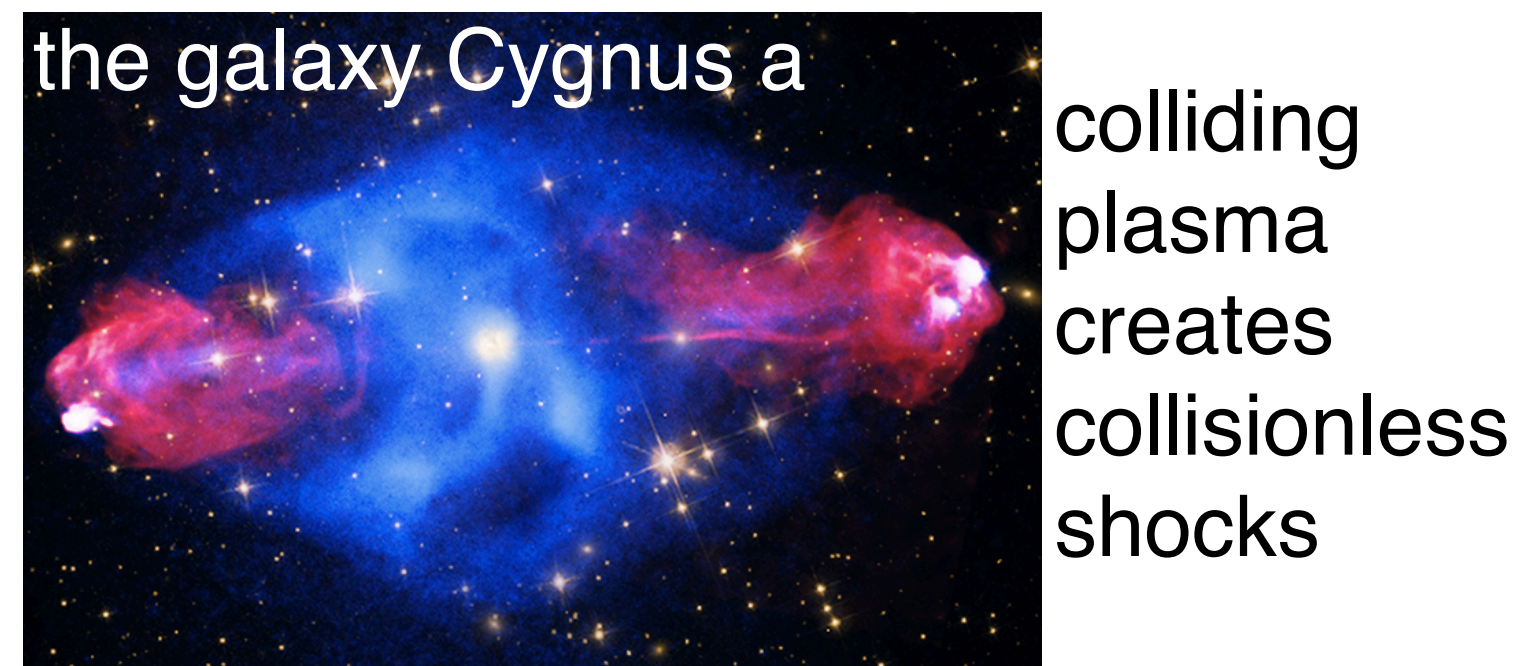
Fast ignition



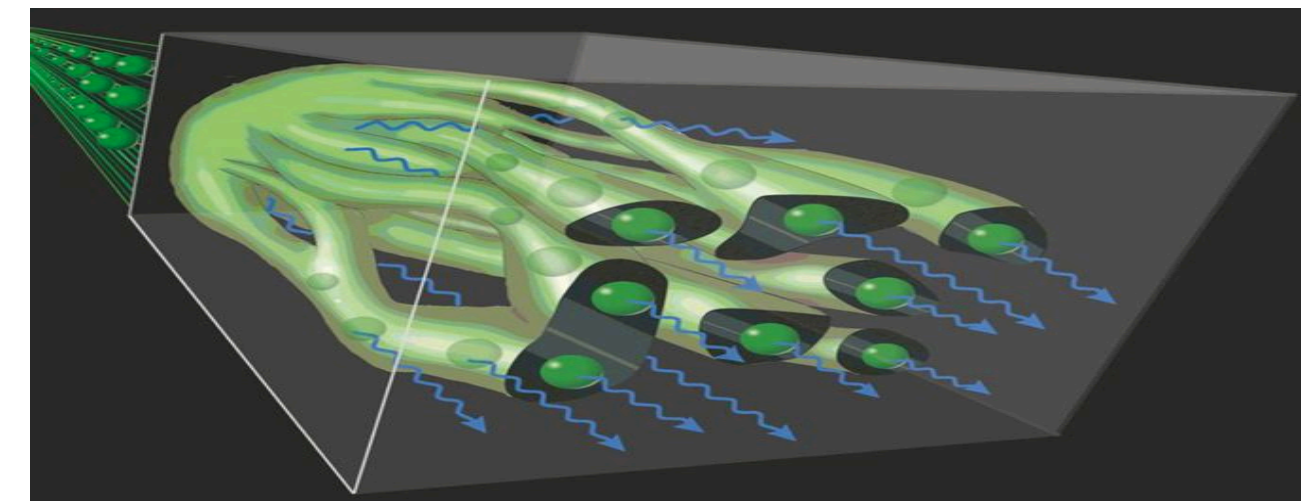
Solar radio burst



Collisionless shock



Gamma ray flash



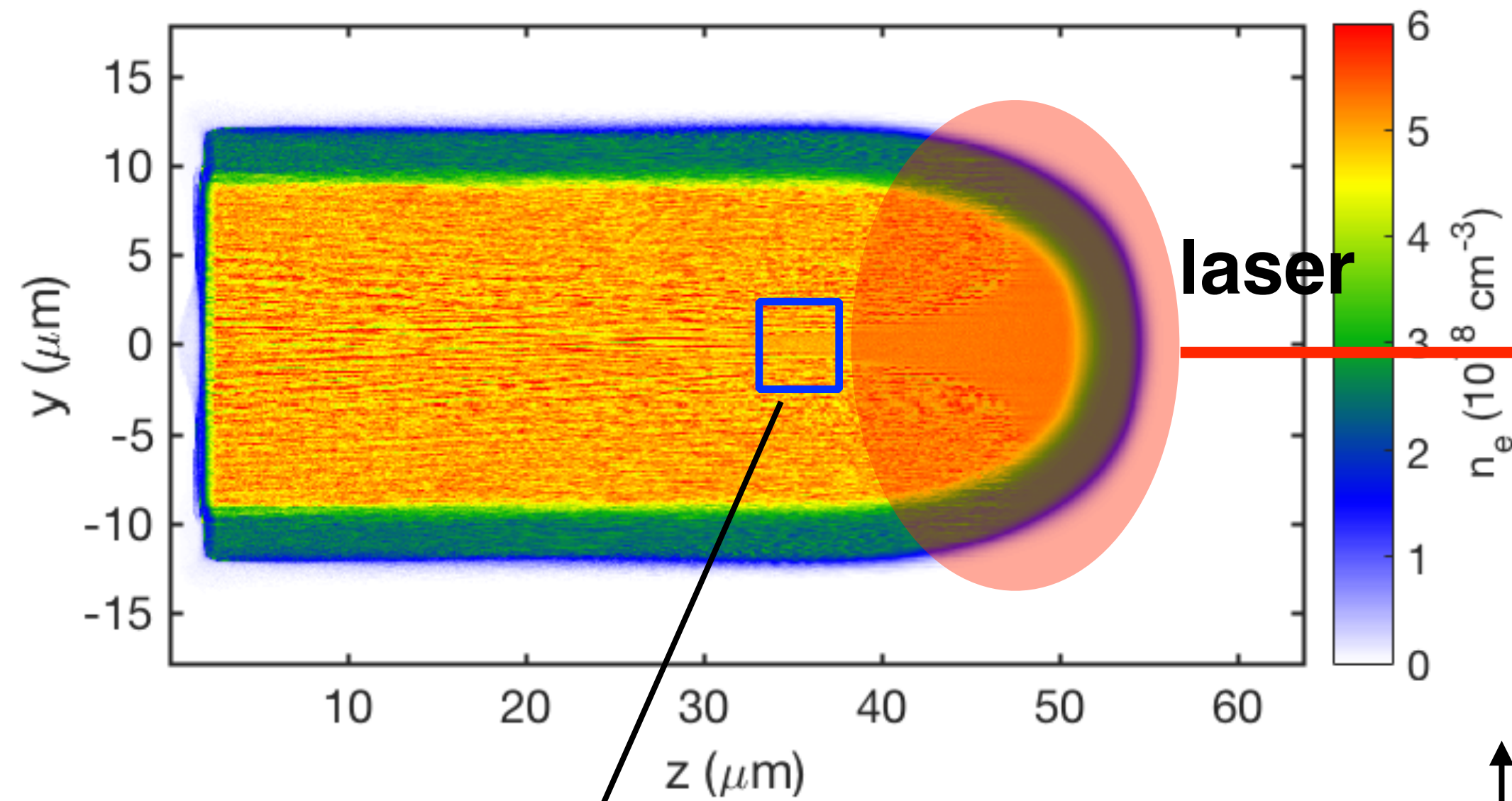
filamentation of e- beam in dense plasma generates bright γ -ray flashes

- One way to study electron current filamentation instability is to send an electron beam through a plasma and to observe the filaments of the e- beam once it breaks up.

B. Allen et al., Phys. Rev. Lett. (2012)

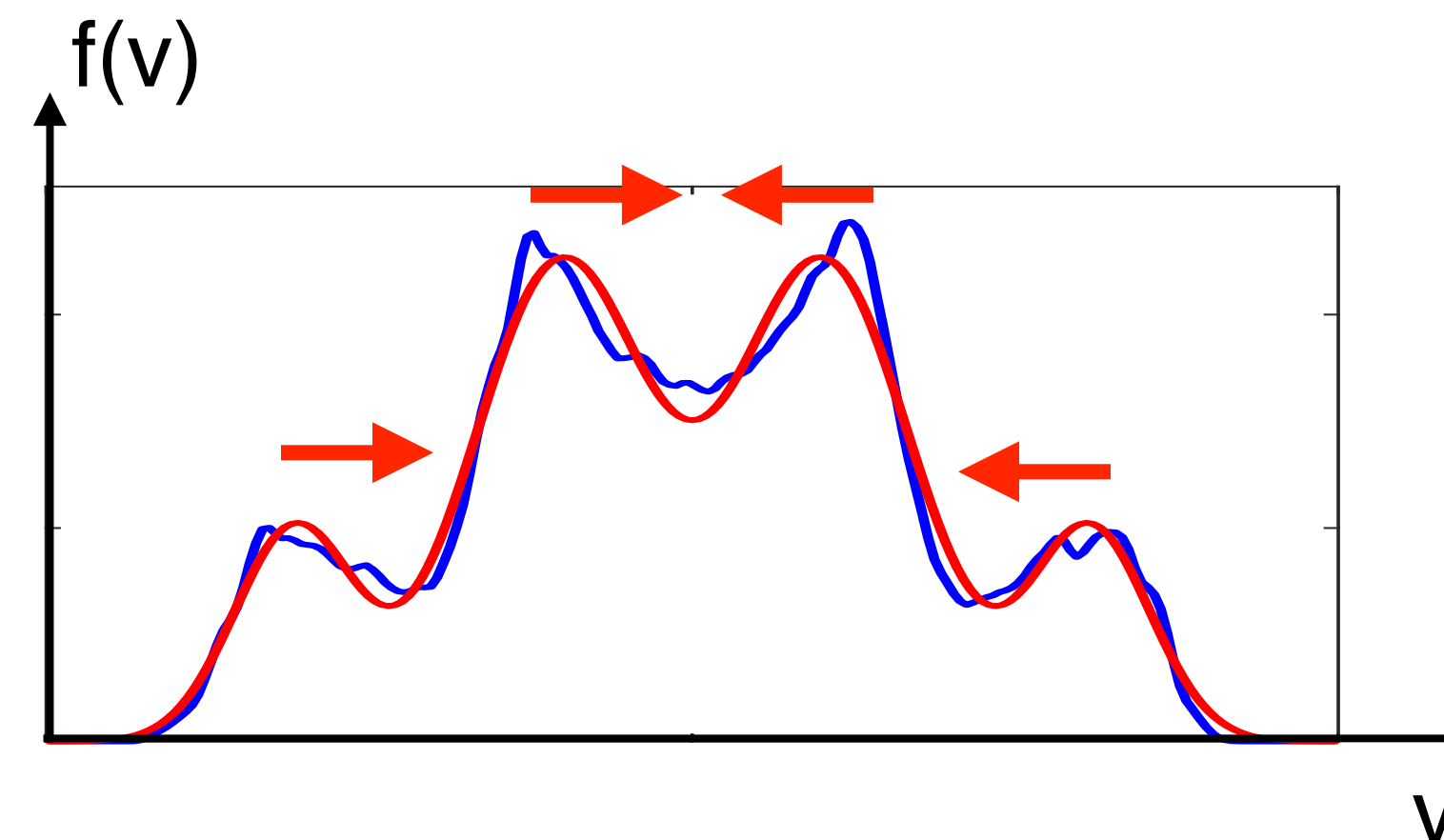
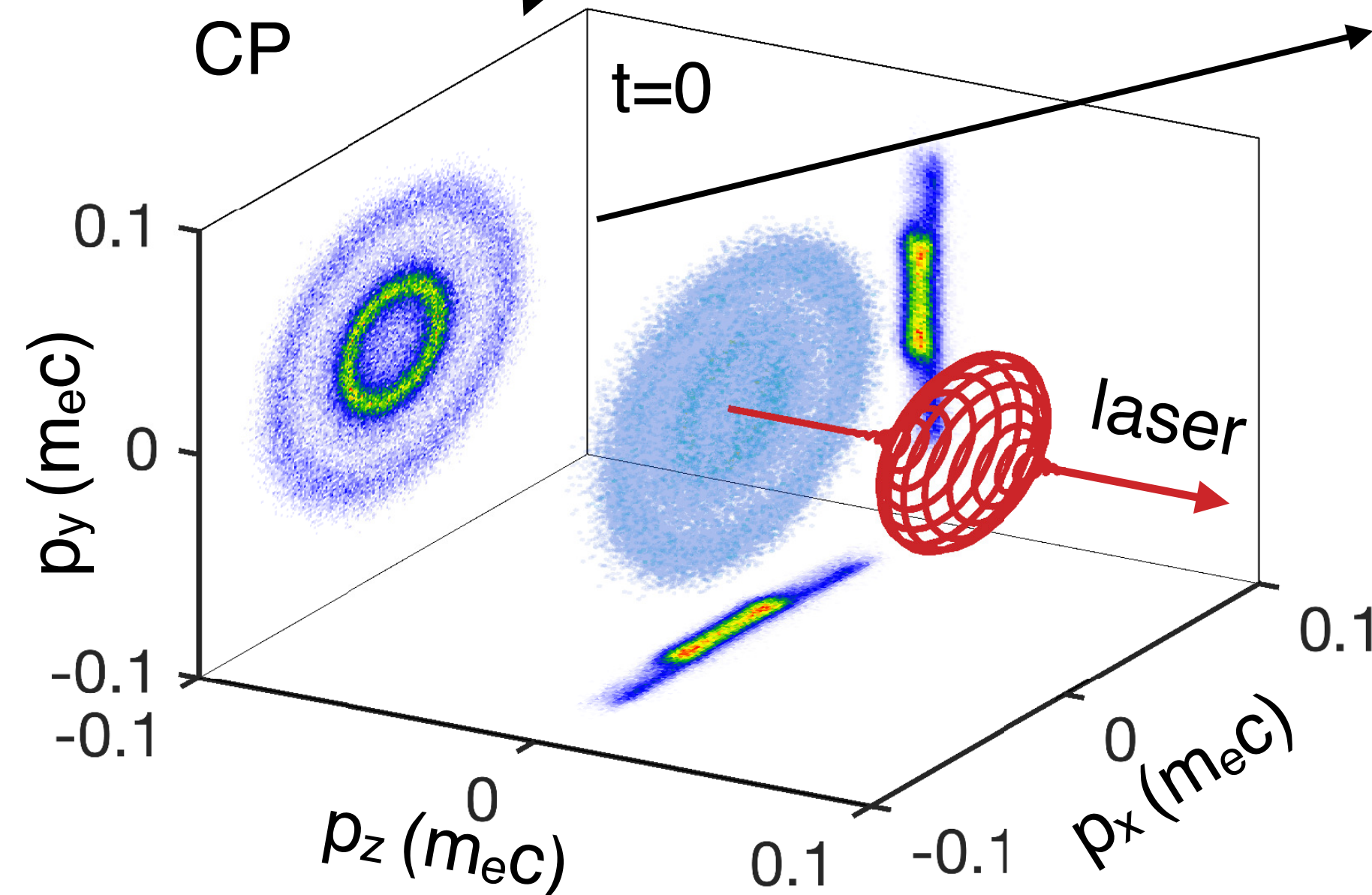
- Recently, we have developed a new laboratory platform for (quantitatively) studying kinetic plasma instabilities.
- We use **optical-field-ionization (OFI)** to initialize plasmas with anisotropic (non-Maxwellian) velocity distributions to trigger kinetic plasma instabilities.
 - streaming instability
 - filamentation instability
 - Weibel instability
- We use **time-resolve Thomson scattering** to probe the plasma density fluctuations associated with these instabilities to get their frequency and the growth rate.

Anisotropic EVD initialized by a circularly polarized laser



laser:
circular polarization
 $\tau=50 \text{ fs}$, $w_0=8 \mu\text{m}$
 $I=1.6 \times 10^{17} \text{ W/cm}^2$
 $a_0 \sim 0.2$

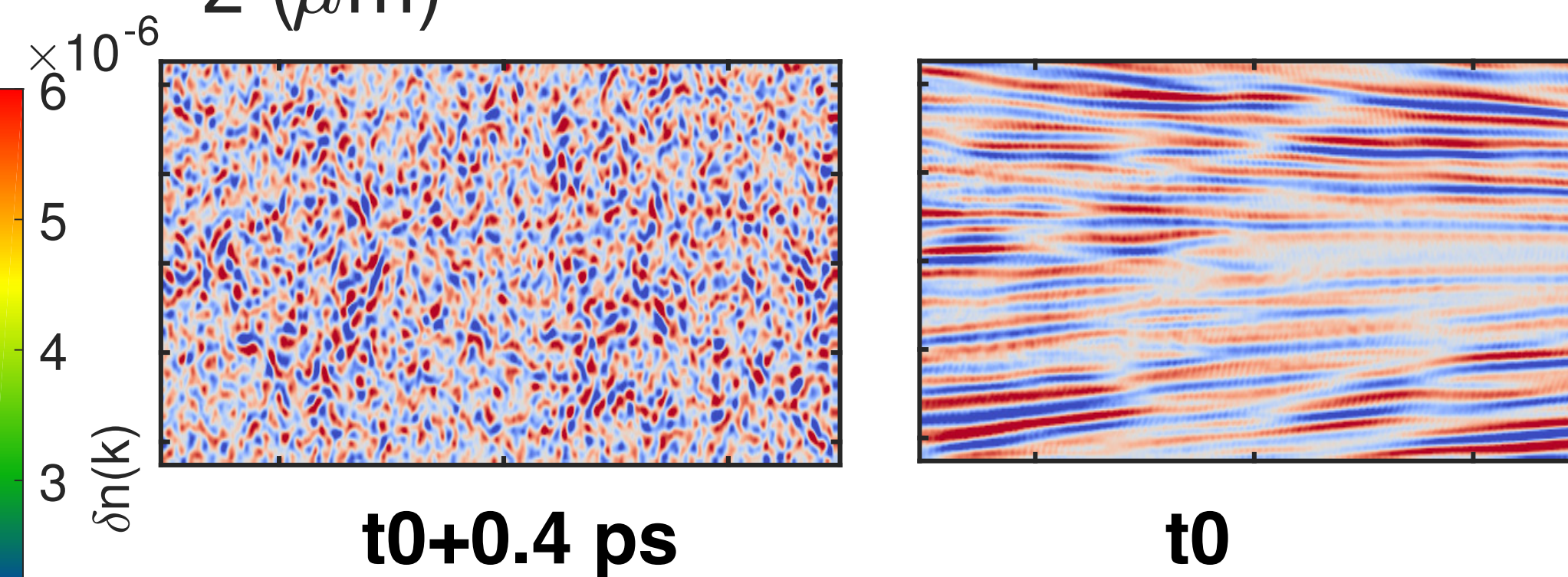
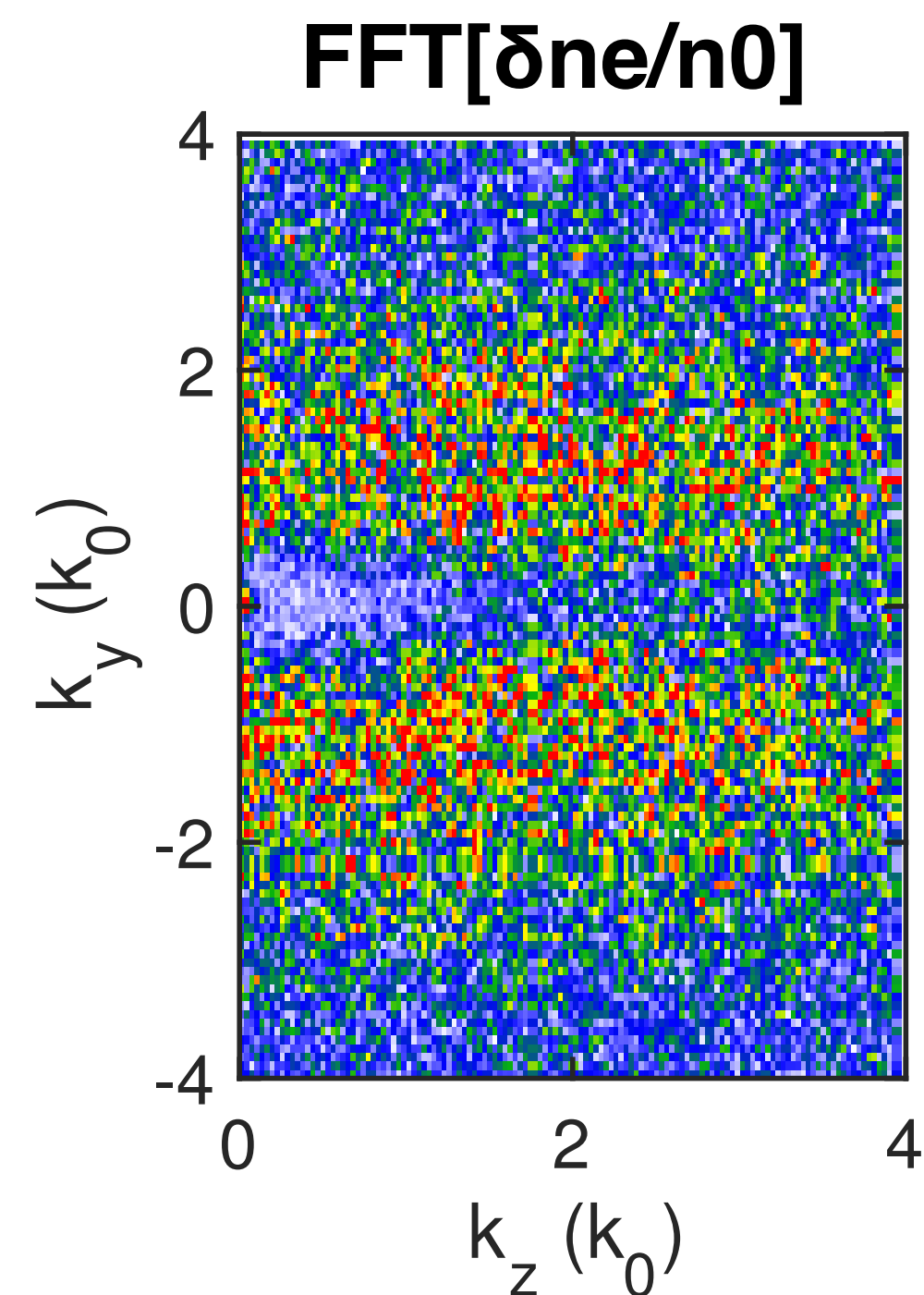
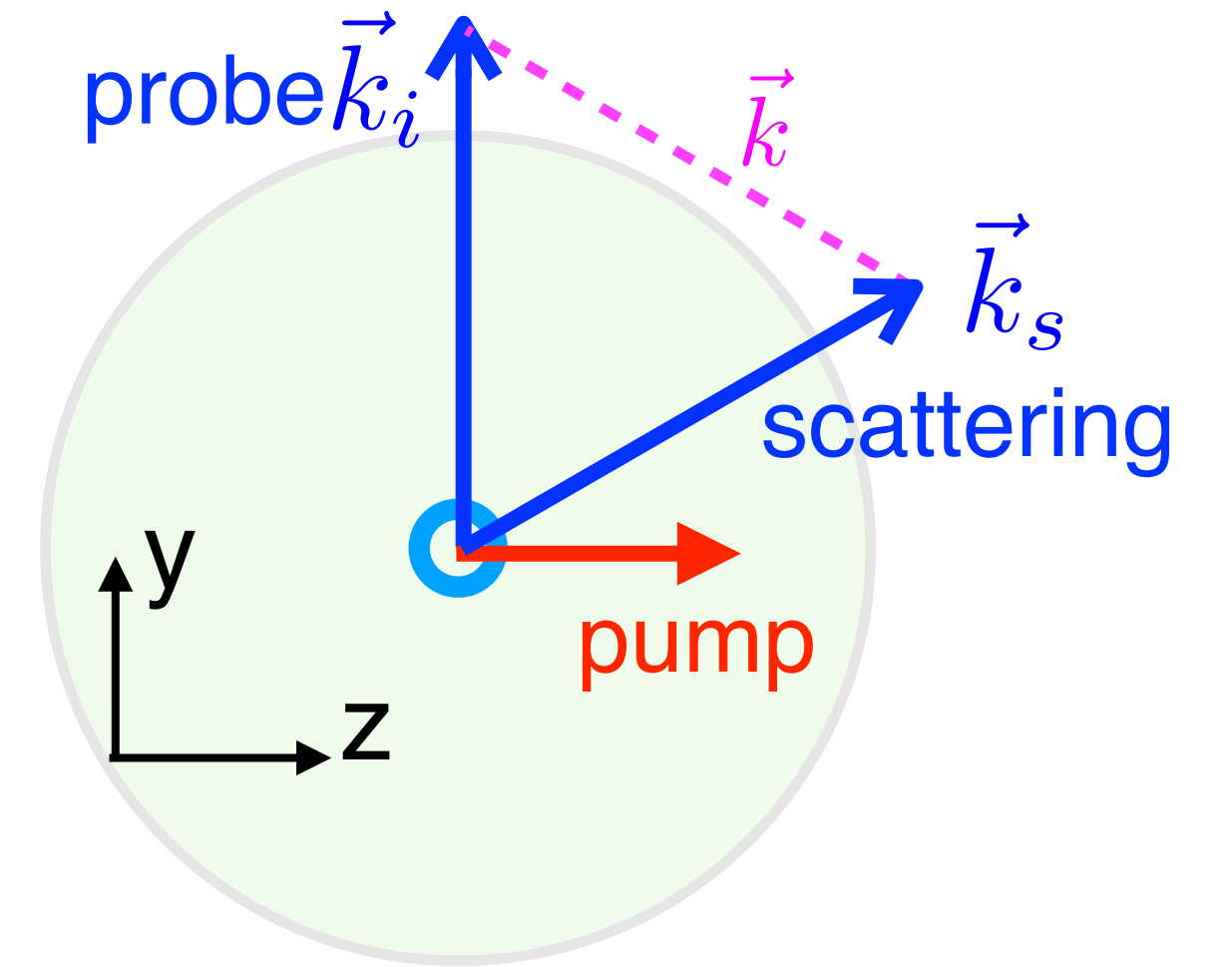
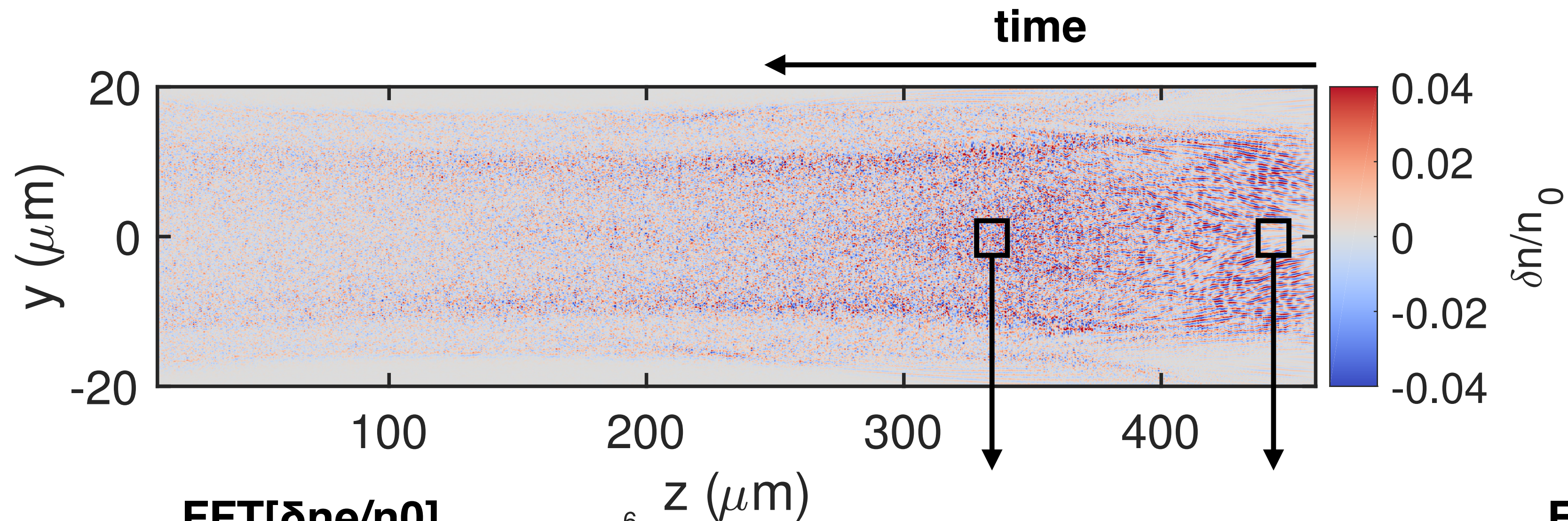
Helium is fully ionized
in $\sim 10 \text{ fs}$, $n_e=5 \times 10^{18} \text{ cm}^{-3}$



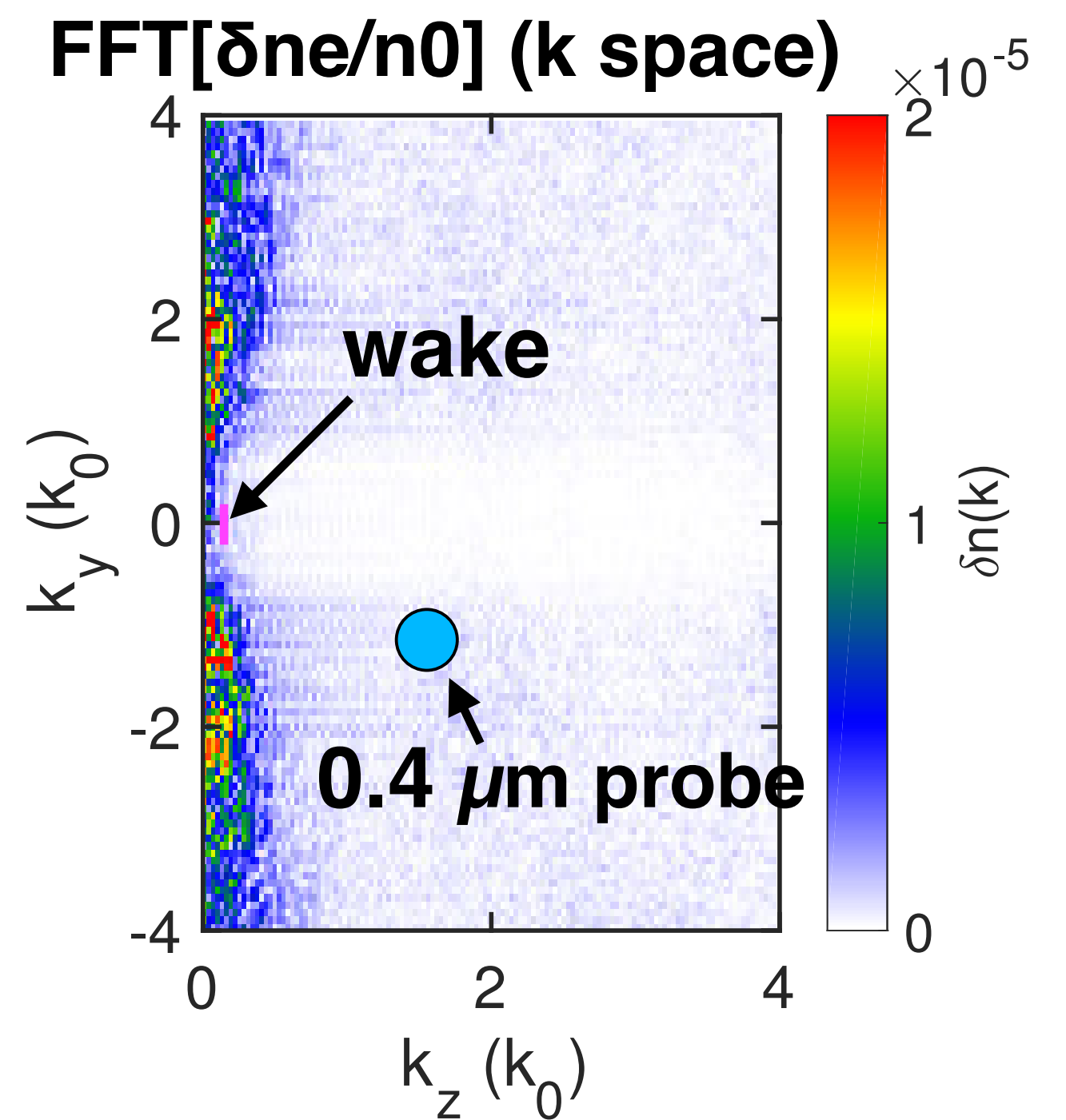
**unstable to streaming
and filamentation
instabilities;**

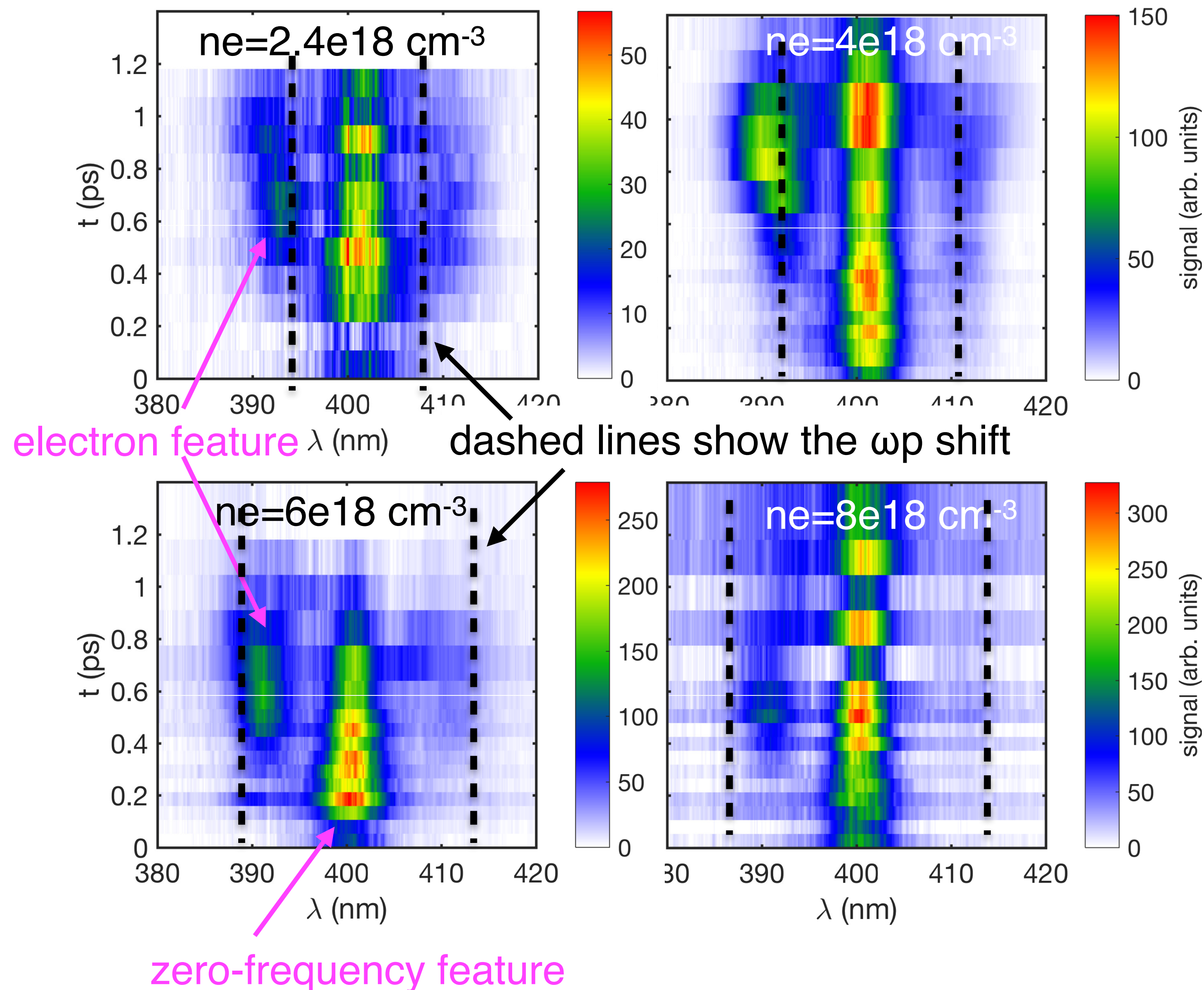
the initial EVD can be
well approximated by four
drifting Maxwellians
(interpenetrating streams)

Signature of instabilities: density fluctuations



The magnitude of density fluctuation is few percent, but is measurable using Thomson scattering





- Laser parameters:
 - pump: $0.8 \mu\text{m}$, 10mJ, 45fs
 - probe: $0.4 \mu\text{m}$, 45fs
 - t_0 accuracy: ~ 100 fs;
 - temporal resolution: ~ 45 fs

Electron feature:

1. shot-lived (compared to collision/recombination time)
2. constant spectral shift

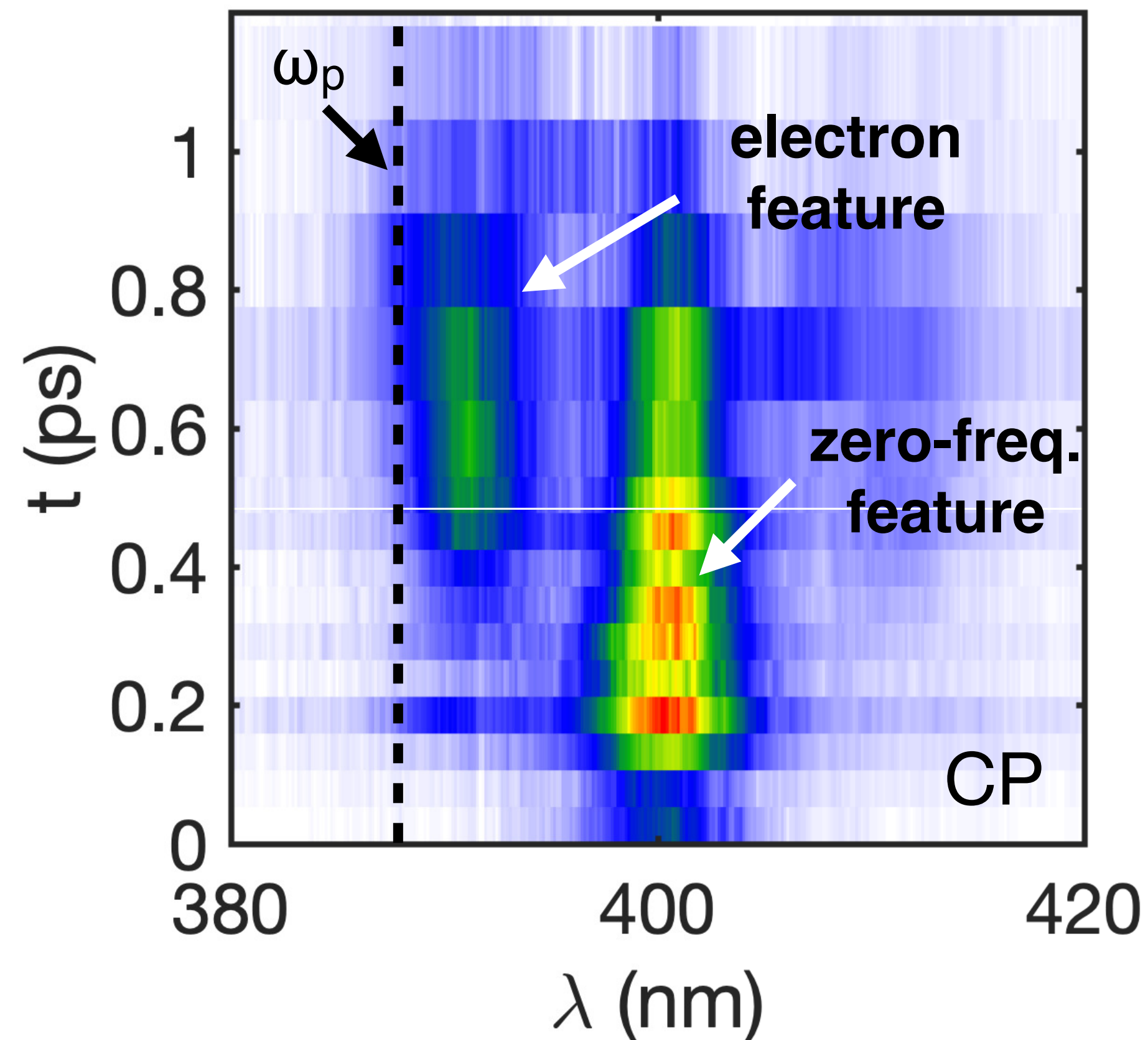
streaming instability

Zero-frequency feature:

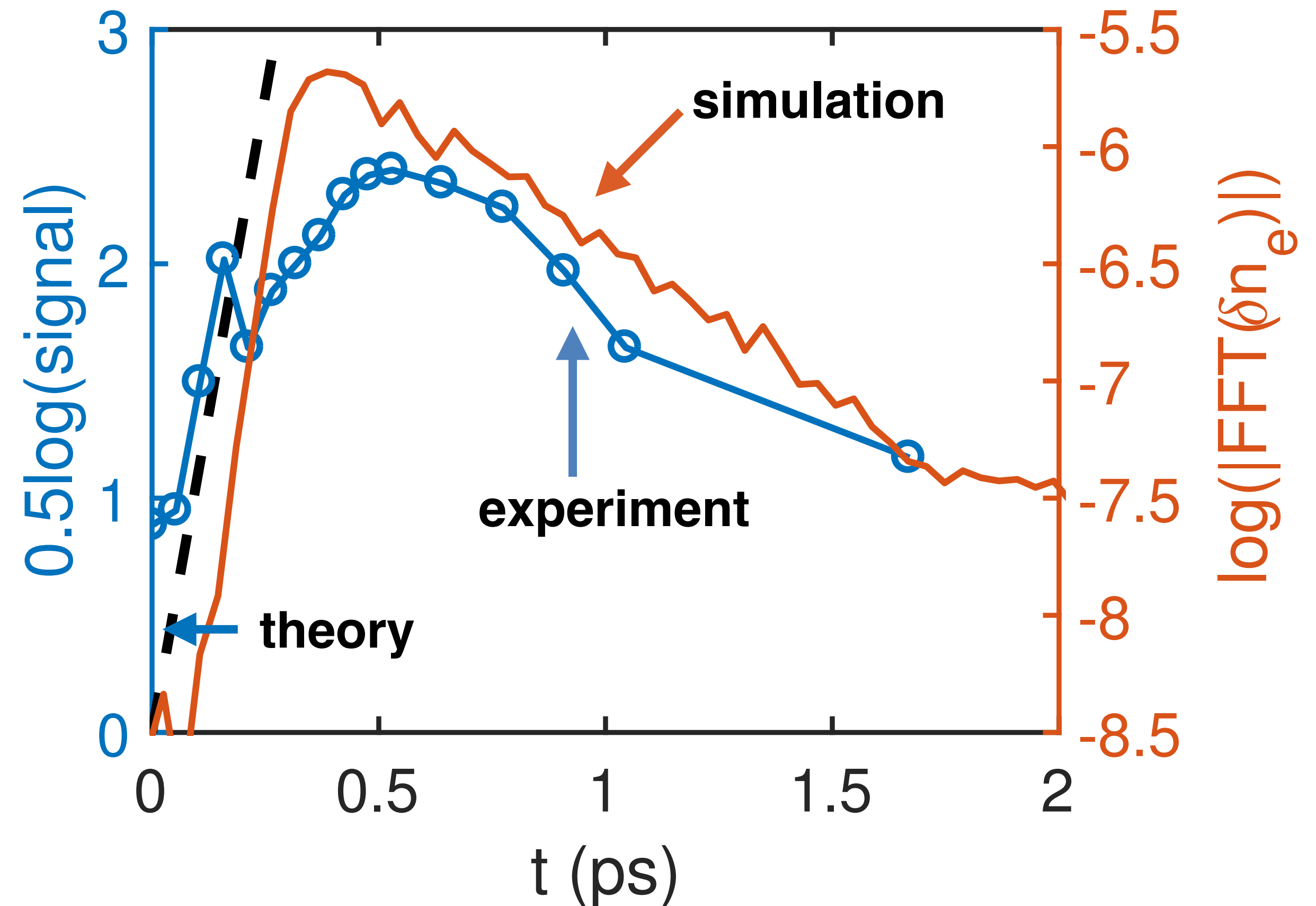
We call it zero-frequency feature instead of “ion feature” because the unshifted signal at $0.4 \mu\text{m}$ also corresponds to instabilities.

filamentation/Weibel instability

measured TS spectrum as a function of time

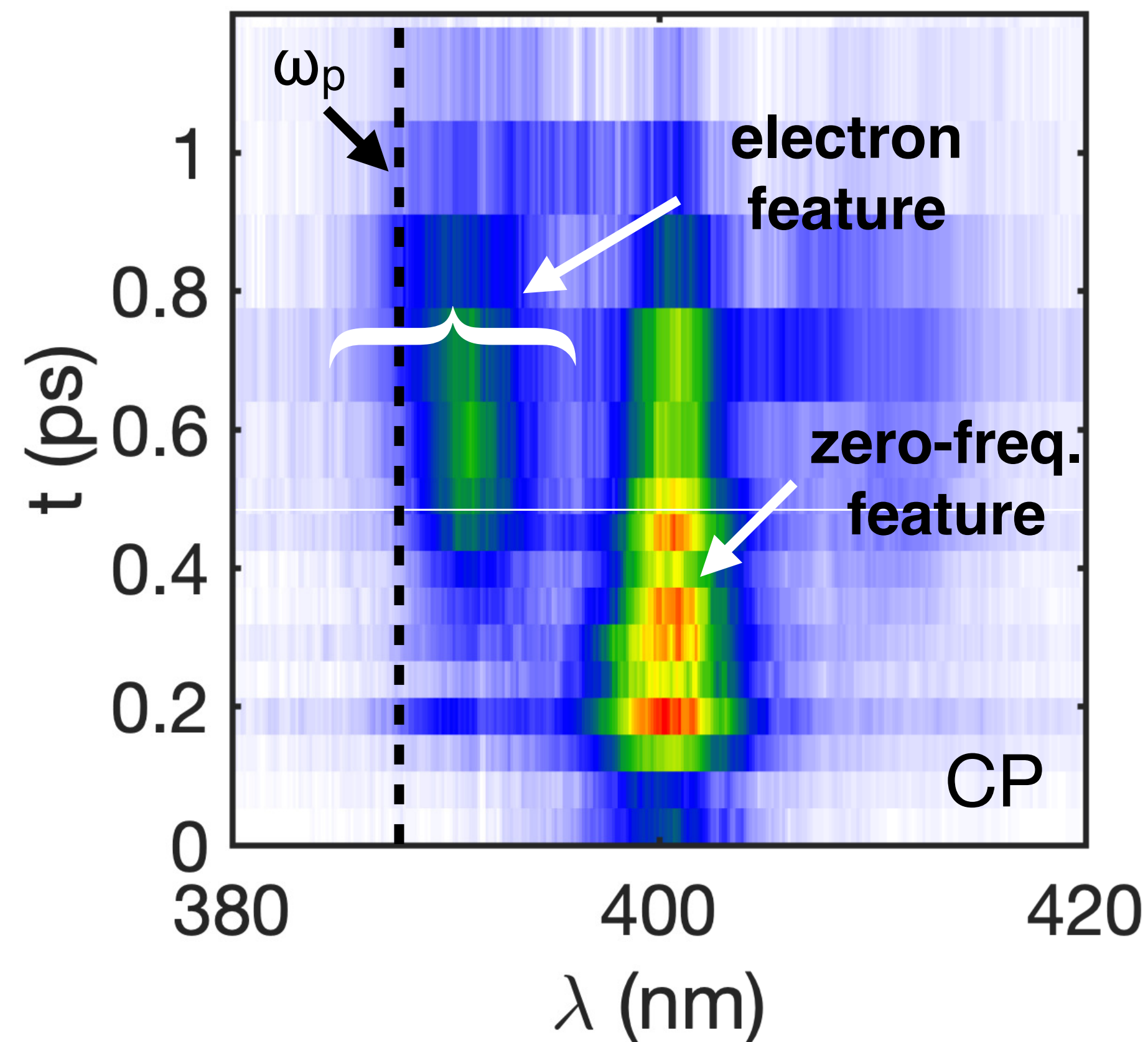


magnitude of the electron feature (streaming instability)

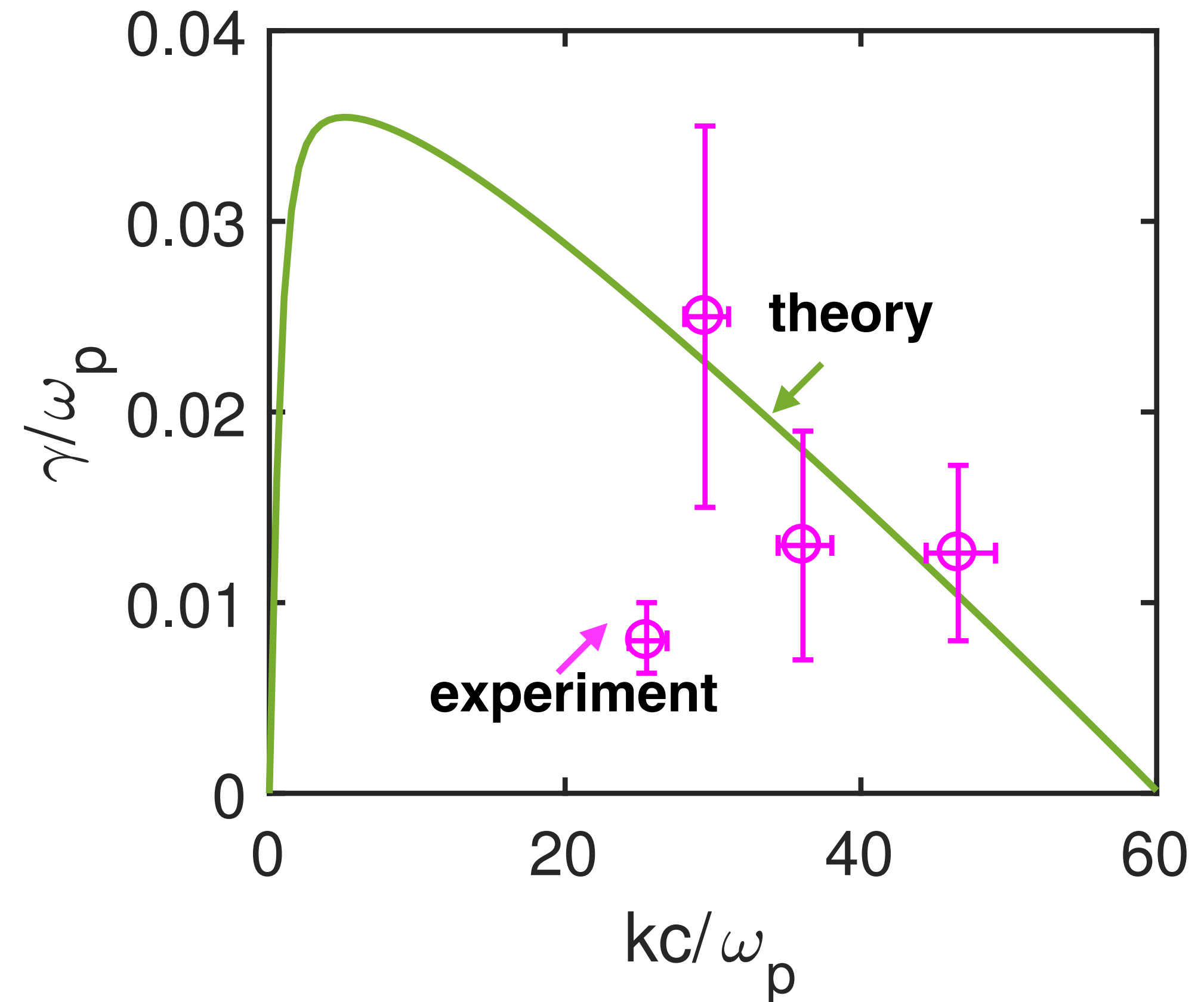
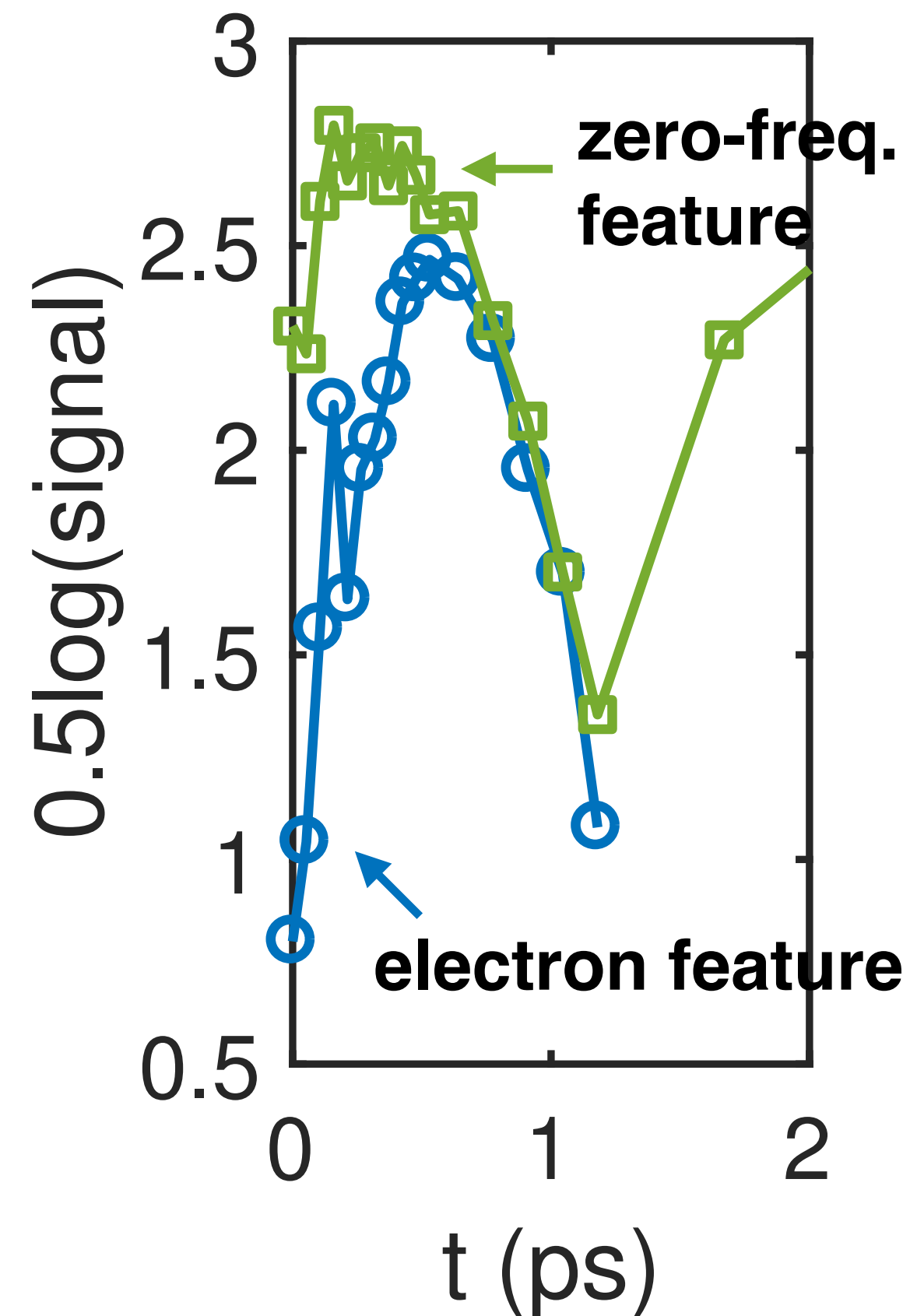


- The linear growth and the nonlinear evolution of the streaming instability were measured;
- Experiment, Kinetic theory and PIC simulation agree well with each other;

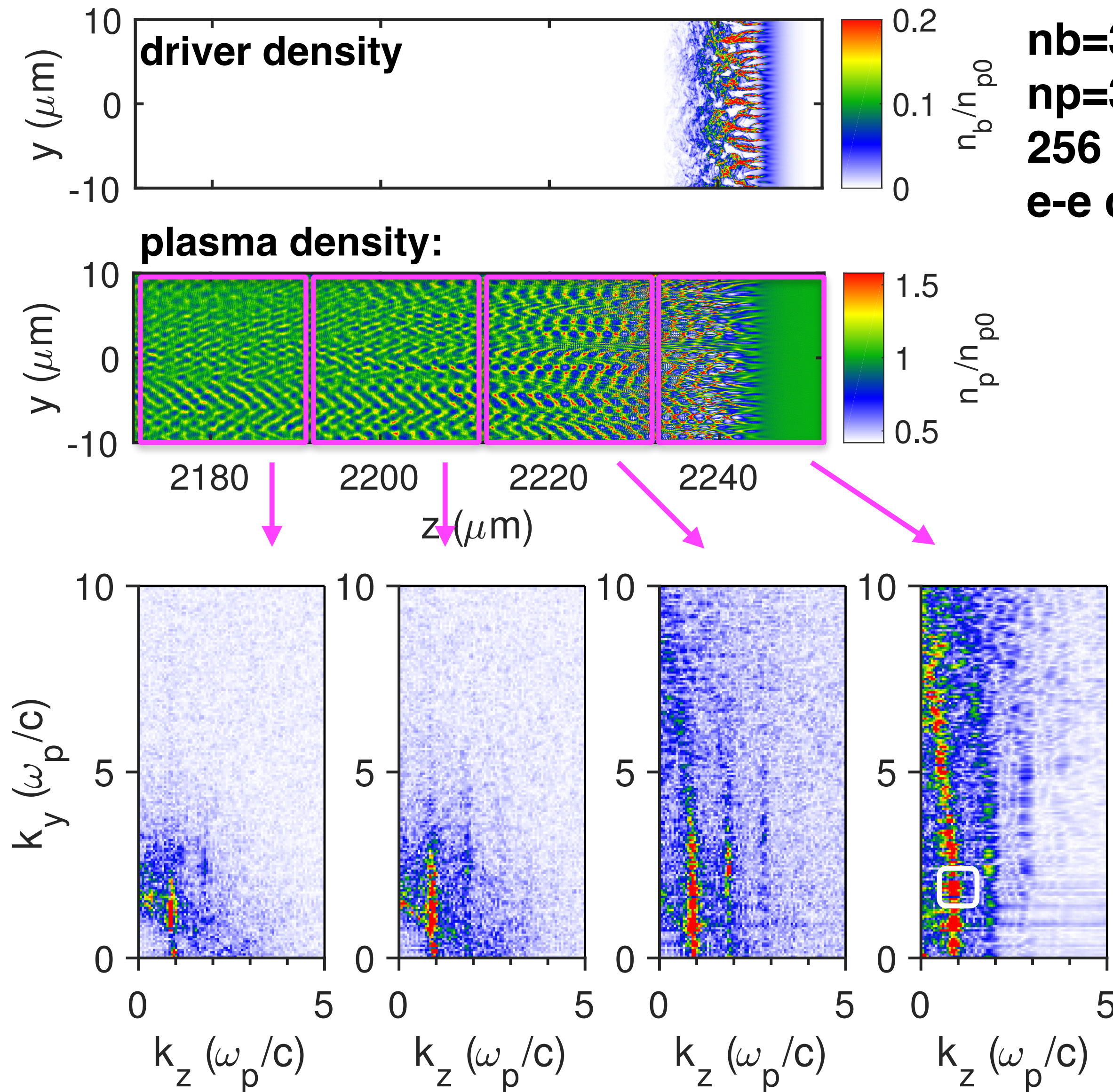
measured TS spectrum as a function of time



The very similar behavior of the electron feature and the zero-frequency feature within the 1 ps suggests that they may have the same driven source- the counter-propagating streams.



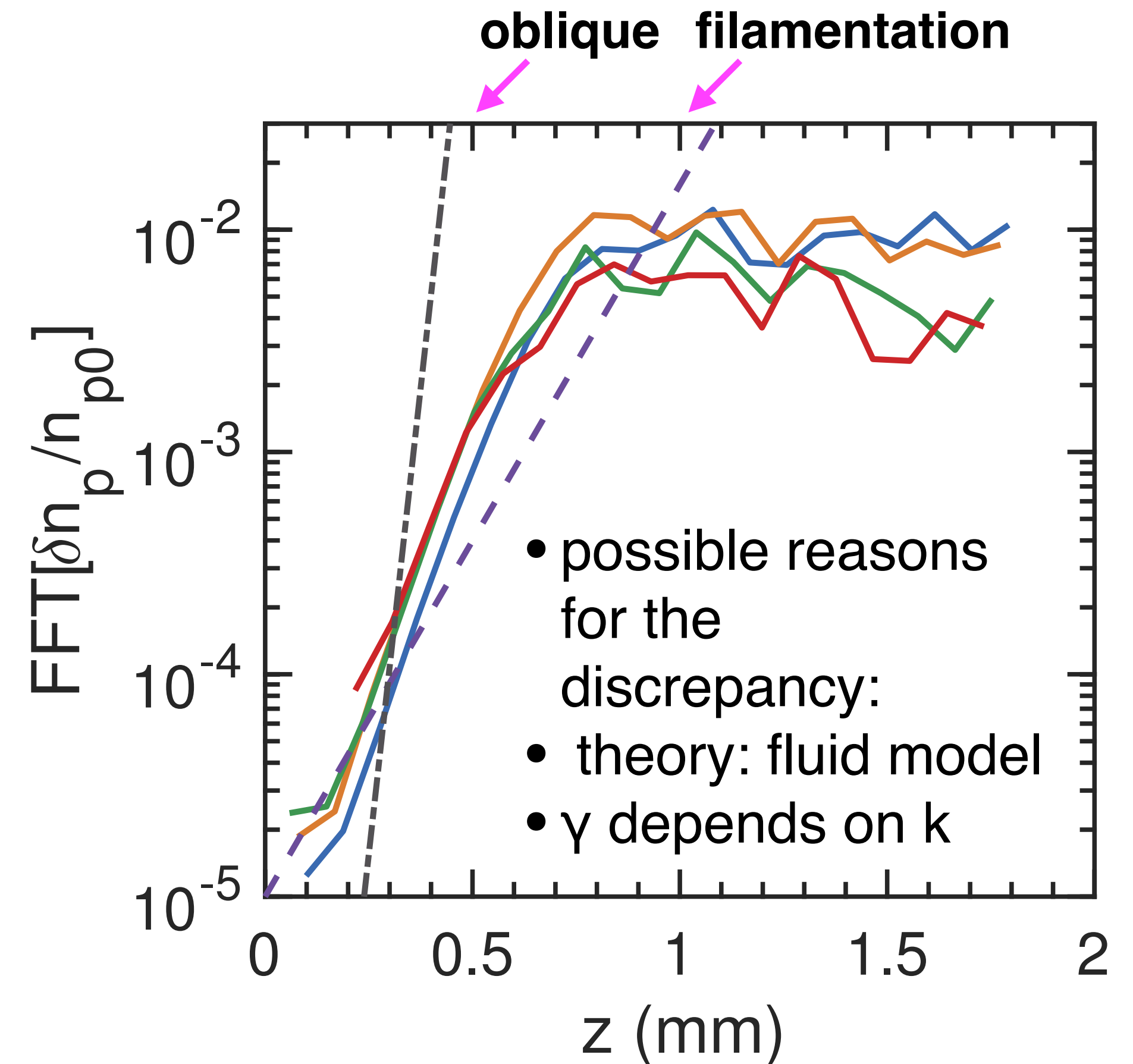
Current filamentation instability: proposed experiment at FACET-II



$n_b=3e19 \text{ cm}^{-3}$, $\sigma_z=2.8 \mu\text{m}$
 $n_p=3e20 \text{ cm}^{-3}$, 20 eV
256 particles/cell
e-e collisions included

- In collaboration with S. Corde at LOA, F. Fiuza, V. Yakimenko, M. Hogan, etc at SLAC

density fluctuation at $k_z=k_p$, $k_y=2k_p$:



- **Theory and PIC simulations show that it is possible to generate high brightness, low emittance e- beams using downramp trapping**
 - **such injected beams may serve as drivers for beam acc. in crystals and nano-structures**
- **Using time-resolved Thomson scattering, we have measured instabilities growing in an OFI plasma to test kinetic theory**
 - **streaming instability**
 - **filamentation instability**
- **Beam current filamentation instability may play important roles in acceleration in Crystals/Nano-structures**